

Calculating ultrafast laser-induced magnetization dynamics

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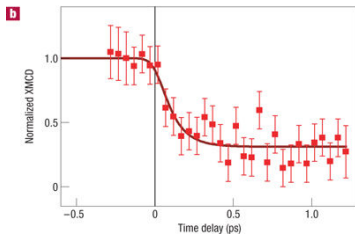
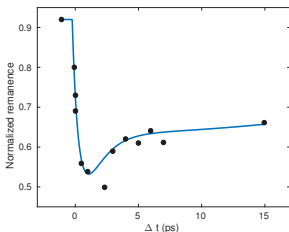
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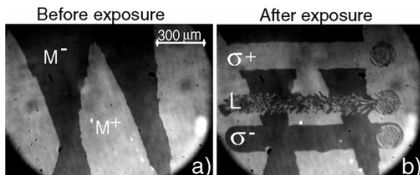
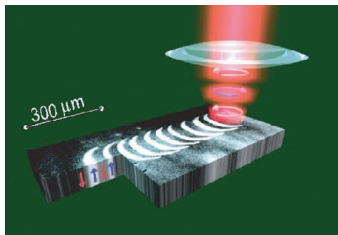
Femtosecond magnetization dynamics

- Magnetization change without the external magnetic field, caused by an intense laser pulse
- *Beaurepaire, 1996*: A significant decrease of magnetization in Ni after an intense laser pulse. (Later observation by Stamm et al.)



- Previously unexplored timescale for magnetism: below 1 ps

Optically controlled magnetization switching

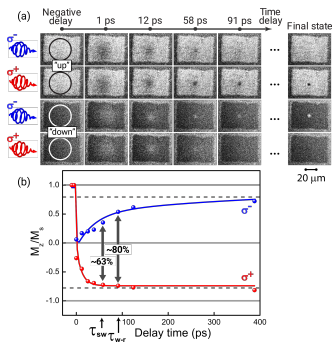


- Stanciu et al., PRL 99 (2007)
- Vahaplar et al., PRL 103 (2009)
- Applications with high impact in data storage - permanently stored information at the speed of RAM
- Origin: Inverse Faraday effect?

Outline

- 1 Magnetization reversal in rare-earth based ferrimagnets
- 2 Demagnetized state in Co

Evolution after the pulse disappears



- $Gd_{24}Fe_{66}Co_{10}$: chemically disordered, amorphous. Close to $GdFe_2$.
- After 1ps: no information about the direction of evolution
- No precession as in field-driven switching

A narrow window with controlled reversal

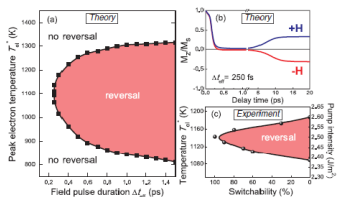


FIG. 2 (color). (a) Phase diagram showing the magnetic state of the $(30 \text{ nm})^3$ volume achieved within 10 ps after the action of the optomagnetic pulse with parameters $H_{eff} = 20 \text{ T}$, $\Delta t_{eff} = 10 \text{ ps}$, and $T_{el}^* = 1130 \text{ K}$. (b) The averaged z component of the magnetization versus delay time as calculated for 250 fs magnetic field pulses $H_{eff} = \pm 20 \text{ T}$ and $T_{el}^* = 1130 \text{ K}$. (c) Switchability versus the pump

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Ultrafast Path for Optical Magnetization Reversal via a Strongly Nonequilibrium State

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Using time-resolved single-shot pump-probe microscopy we unveil the mechanism and the time scale of all-optical magnetization reversal by a single circularly polarized 100 fs laser pulse. We demonstrate that the reversal has a linear character, i.e., does not involve precession but occurs via a strongly nonequilibrium state. Calculations show that the reversal time which can be achieved via this mechanism is within 10 ps for a 30 nm domain. Using two single subpicosecond laser pulses we demonstrate that for a 5 μm domain the magnetic information can be recorded and readout within 30 ps, which is the fastest "write-read" event demonstrated for magnetic recording so far.

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Reversal induced by each pulse



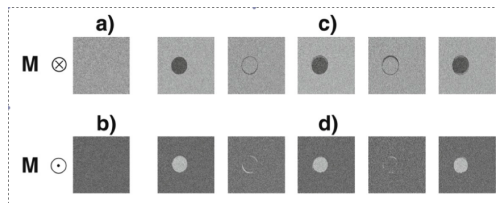
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Ultrafast heating as a sufficient stimulus for magnetization reversal in a ferrimagnet

T.A. Costler¹, J. Barker¹, R.F.L. Evans¹, R.W. Chantrell¹, U. Atxitia², O. Chubykalo-Fesenko², S.E. Moussaoui³, L. Le Guyader³, E. Mengotti³, L.J. Heyderman³, F. Nolting³, A. Tsukamoto⁴, A. Ichik⁵, D. Afanasiev⁵, B.A. Ivanov⁵, A.M. Kalashnikova⁶, K. Vahaplar⁷, J. Mentink⁷, A. Kirilyuk⁷, Th. Rasing⁷ & A.V. Kimel⁷

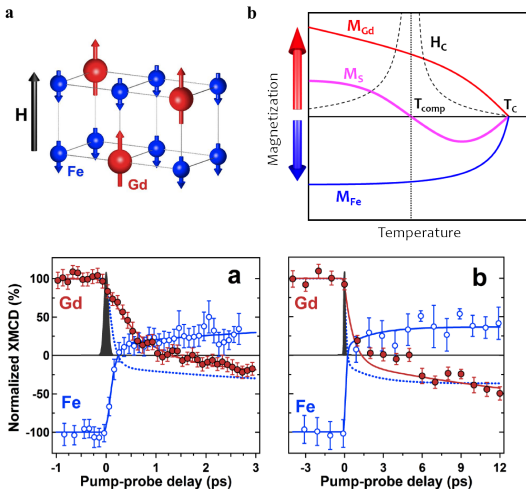


Discrepancy in experiments?

- Inverse Faraday effect ruled out, the original polarization dependence was only due to different absorption coefficients for different polarizations (happening in a very narrow fluence window)
- No external field, circular polarization of the laser pulse or any other direct source of angular momentum, but: magnetization acquires specific direction different from the original one

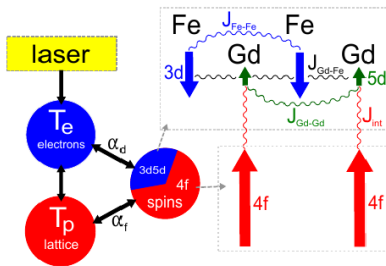
Time-resolved XMCD study

Radu et al., Nature 472 (2011) :



Model (co-operation with Konstanz)

- GdFe_2 : cubic Laves phase (C_{15}), $T_C=560\text{K}$, $T_M=300\text{K}$
- Gd: 90% of mag. moment in 4f states, but localized more than 4eV below E_F
- $J_{TM-TM} = 32.5 \text{ meV}$, $J_{R-R} = 7.8 \text{ meV}$, $J_{TM-R} = -3.2 \text{ meV}$,
 $J_{int} = 130 \text{ meV}$



Simulations based on LLG equation

$$\mathcal{H} = - \sum_{\langle ij \rangle} \frac{J_{ij}}{2} \mathbf{S}_i \cdot \mathbf{S}_j - \sum_{i \in \text{Gd}} J_{int} \mathbf{S}_i \cdot \mathbf{S}'_i - d_z \sum_i (S_i^z)^2. \quad (1)$$

Landau-Lifschitz-Gilbert equation:

$$\frac{\partial \mathbf{S}_i}{\partial t} = - \frac{\gamma_i}{(1 + \alpha_i^2) \mu_S^i} \mathbf{S}_i \times \mathbf{H}_i(t) - \frac{\alpha_i \gamma_i}{(1 + \alpha_i^2) \mu_S^i} \mathbf{S}_i \times (\mathbf{S}_i \times \mathbf{H}_i(t)).$$

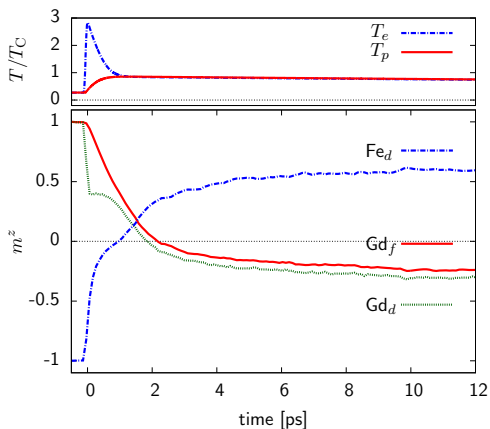
γ_i ... the gyromagnetic ratio. For the 4f electrons of Gd: the same equation with the primed quantities.

Thermal fluctuations: additional white-noise term - the effective field

$\mathbf{H}_i(t) = - \frac{\partial \mathcal{H}}{\partial \mathbf{S}_i} + \zeta_i(t)$, the thermal noise term ζ_i

$$\langle \zeta_i(t) \rangle = 0, \quad \langle \zeta_{i\eta}(0) \zeta_{j\theta}(t) \rangle = \delta_{ij} \delta_{\eta\theta} \delta(t) 2\alpha_i k_B T_i \mu_S^i / \gamma_i. \quad (2)$$

Simulated magnetization evolution



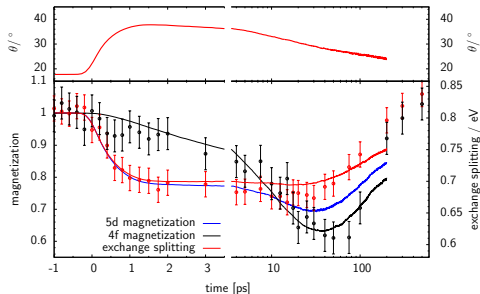
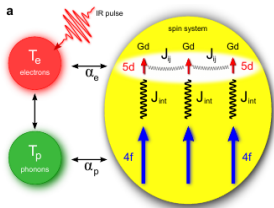
S. Wienholdt, D. Hinzke, K. Carva, P. M. Oppeneer, and U. Nowak, *Phys. Rev. B* 88, 020406(R) (2013)

Gd: LLG with 5d momentum calculation

- Ab initio exchange constants employed in LLG simulation
 (B. Frietsch, KC, PMO et al. , Nat Commun 6, 2015):

$$\frac{\partial \mathbf{S}_i}{\partial t} = - \frac{\gamma_i}{(1+\alpha_i^2)\mu_B^i} \mathbf{S}_i \times \mathbf{H}_i(t) - \frac{\alpha_i \gamma_i}{(1+\alpha_i^2)\mu_B^i} \mathbf{S}_i \times (\mathbf{S}_i \times \mathbf{H}_i(t)).$$

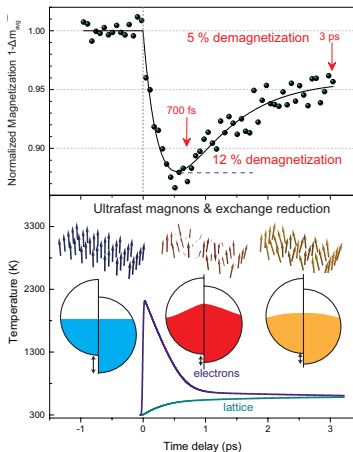
- Mean M, ϑ from previous LLG time step \rightarrow 5d momentum recalculated



Magnetic excitations in a fs laser pumped system

Three possible effects:

- Disorder of magnetic moments (transversal spin excitation)
- Reduction of exchange splitting (longitudinal spin excitation)
- An increase of electronic temperature



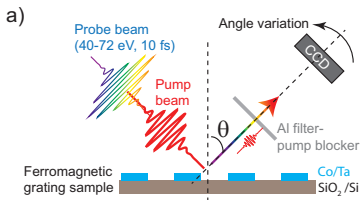
MO response

- High harmonics used to generate EUV light and probe the excited system
MO response (*Fan, DL, KC, PMO, et al., PNAS 112 (2015) 14206*)
- MO asymmetry

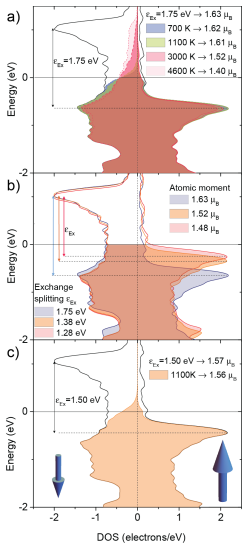
$$A = 2\text{Re} \frac{\epsilon_{xy} \sin 2\vartheta_i}{n^4 \cos \vartheta_i^2 - n^2 + \sin \vartheta_i^2}$$

measured for different energies and angles

- Comparison between the calculated and measured asymmetry spectra
(*Turgut, DL, KC, PMO, et al., Phys. Rev. B 94 (2016) 220408*)



Calculated Co DOSes for different excitations



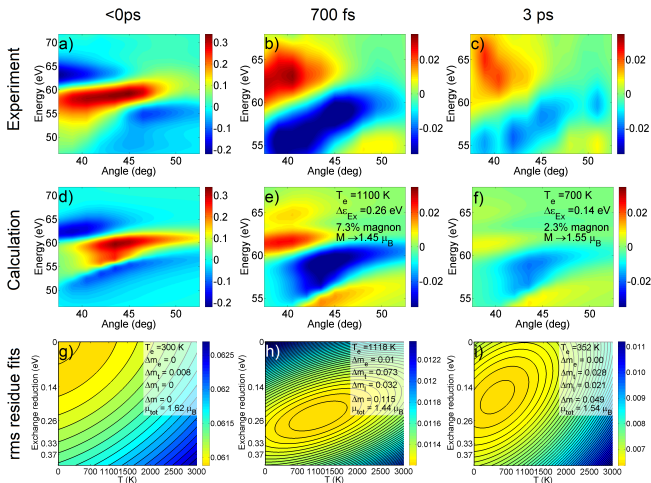
a) Elevated T_e : multiplication with a FD distribution

b) Reduced exchange: fixed spin moment calculation

c) Combination of the above

(Wien2K FP-LAPW code: 1.5mil. corehours,
 $R_{kmax}=8.5 \text{ Ry}^{-1}$, $l_{max} = 10$, kgrid 44x44x22)

Extraction of the contributions



- Compare *Eich et al., 2017*: only magnons found, a conducting substrate

Conclusions

- Ultrafast heating found to be sufficient stimulus for magnetization reversal - the reversal process is also of thermal origin similar to demagnetization
- Breaking of intra-atomic $4f$ - $5d$ coupling in Gd on a ps timescale (*Frietsch et al., Nat Commun, 2015*)
- TR-MO asymmetry measurement allows us to disentangle the different contributions to demagnetization (*Turgut, DL, KC, PMO, et al., PRB 94, 2016*)
- Both transverse and longitudinal excitations found to contribute, at 700fs $\Delta m_t / \Delta m_l > 2$, later the contribution from longitudinal component is growing

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