

Transfer of heat and liquids through the high-pressure ice layer of Ganymede

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Outline



1. Ocean worlds
2. Model & numerical implementation
3. Results
4. Conclusions & perspectives

Ocean worlds

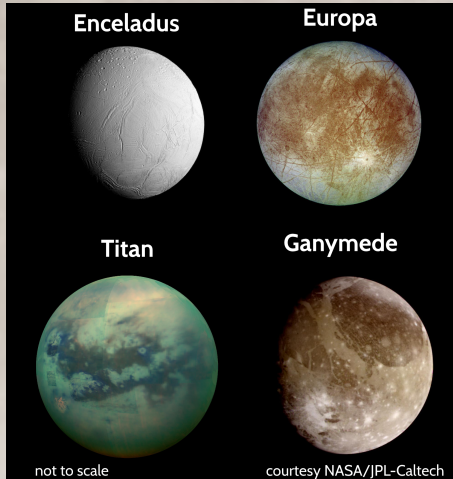
- ▶ water-rock interactions → rich communities of organisms



- ▶ leaching of volatiles (Ar, methane, ...) from the silicates

Ocean worlds

- ▶ water-rock interactions → rich communities of organisms
- ▶ Galileo & Cassini: deep oceans within the icy moons → **ocean worlds**



Formation and evolution of Ocean Worlds in the Solar System and beyond



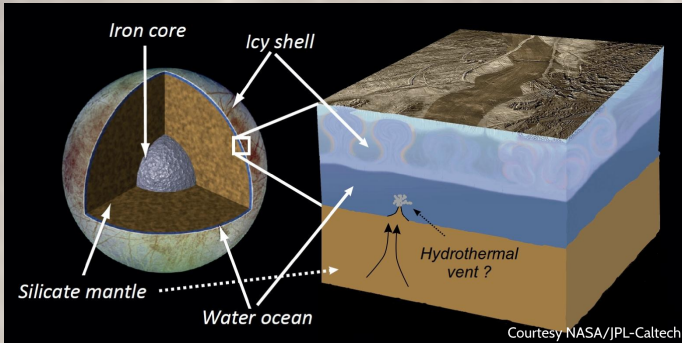
Christophe Sotin
(JPL-Caltech/NASA)

úterý 7. listopadu v 18 hodin
Matematicko-fyzikální fakulta UK
V Holešovičkách 2, posluchárna T1

Vodní světy jsou planety/měsíce, které hostí velké množství kapalné vody na povrchu nebo v nitru. Jejich výzkum je motivován především otázkou možnosti vzniku života mimo Zemi. Ve Sluneční soustavě jsou vodní světy zastoupeny několika ledovými měsíci velkých planet a pozorování mise Kepler přinesla důkazy o existenci mnoha dalších vodních světů i mimo naši soustavu. Přednáška prof. Sotina se zaměří na otázky týkající se formace a evoluce vodních světů. Přednášející, který je jedním z řídicích pracovníků JPL-Caltech/NASA, také představí projekty, které NASA v souvislosti s vodními světy plánuje v příštích letech a rád zodpoví i další dotazy. Profesor Sotin přijíždí na pozvání Katedry geofyziky MFF UK.

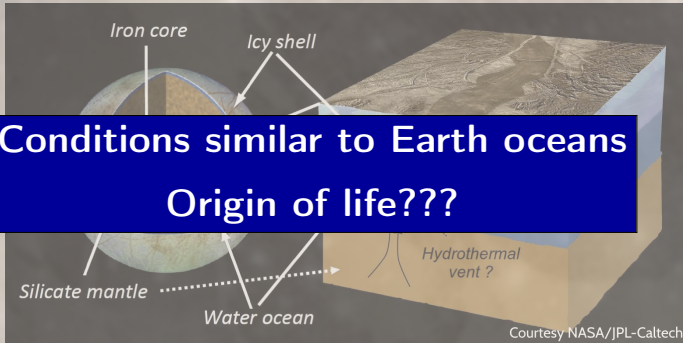
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- ▶ Europa & Enceladus: direct contact of silicates with the ocean



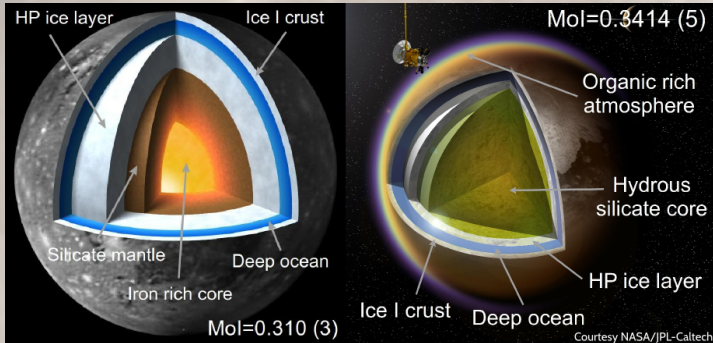
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- ▶ Ganymede & Titan: thick HP ice layer between silicates and ocean



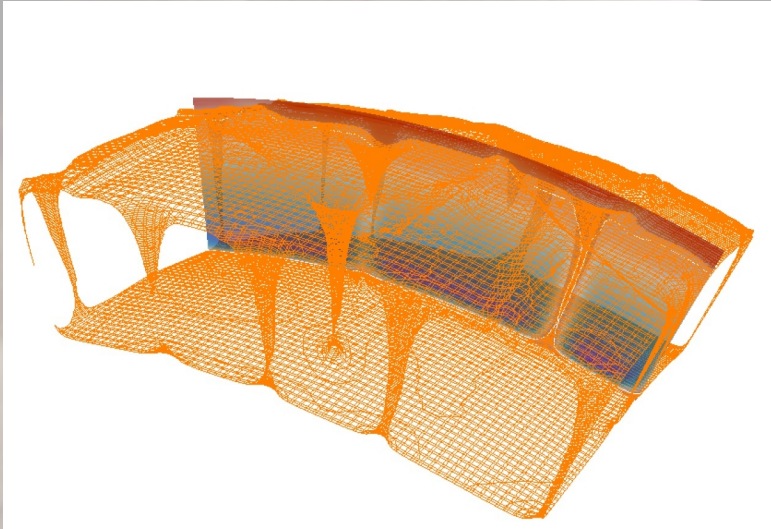
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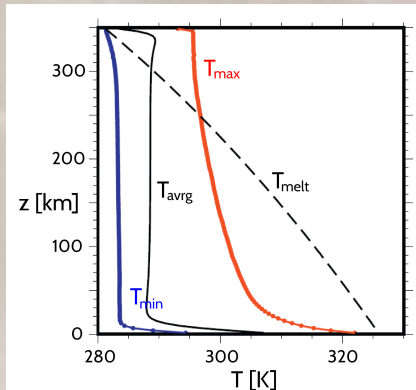
Heat transfer through HP ice layer

- ▶ *Choblet et al. (2017)*: solid-state thermal convection



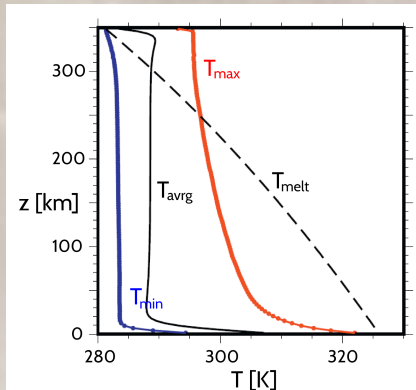
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- cold thermal boundary layer cannot exist in a solid state



Heat transfer through HP ice layer

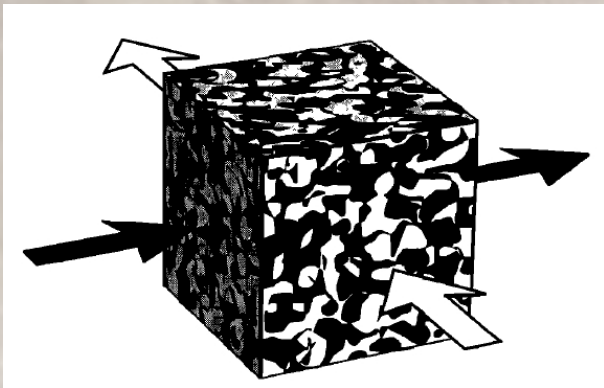
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→ **two-phase mixture model**

Two-phase mixture model

- ▶ coexisting phases: solid ice matrix (i) + liquid water (w)
- ▶ volume fraction of water in the mixture: porosity $\phi = \frac{V_w}{V_i + V_w}$
- ▶ amount of ice: $1 - \phi$
- ▶ cold ice ($\phi = 0$) vs. temperate ice ($\phi \sim$ few percents)
- ▶ assumption of thermal equilibrium in the temperate ice: $T_i = T_w = T_m$



Governing equations

► Stokes system $\nabla \cdot \mathbf{v}_i = \frac{\Delta \rho}{\rho_w \rho_i} r$

$$\begin{aligned} \nabla \Pi = & - (1-\phi) \rho_i \alpha_i (T - T_m) \mathbf{g} - \phi \Delta \rho \mathbf{g} + \nabla \left(\frac{(1-\phi) \mu_i}{\phi} (\nabla \cdot \mathbf{v}_i) \right) \\ & + \nabla \cdot \left((1-\phi) \mu_i \left[\nabla \mathbf{v}_i + (\nabla \mathbf{v}_i)^T - \frac{2}{3} (\nabla \cdot \mathbf{v}_i) \mathbf{I} \right] \right) \end{aligned}$$

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- extension ($r > 0$)
compaction ($r < 0$)

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- extension/
/compaction
- **thermal** &
phase-change
buoyancy

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- ▶ extension/
/compaction
- ▶ thermal &
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buoyancy
- ▶ **bulk** & **shear** viscous
deformation

Governing equations

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- ▶ energy balance

$$\rho_i c_i \frac{\partial T}{\partial t} + \rho_i c_i \mathbf{v}_i \cdot \nabla T + Lr = k_i \nabla^2 T$$

- ▶ extension/
/compaction
- ▶ thermal &
phase-change
buoyancy
- ▶ bulk & shear viscous
deformation
- ▶ heat consumption
($r > 0$) release ($r < 0$)

Governing equations

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▶ advection of porosity

$$\frac{\partial \phi}{\partial t} + \nabla \cdot (\phi \mathbf{v}_w) = \frac{r}{\rho_w}$$

- ▶ extension/
/compaction
- ▶ thermal &
phase-change
buoyancy
- ▶ bulk & shear viscous
deformation
- ▶ heat consumption/
/release
- ▶ melting ($r > 0$)
freezing ($r < 0$)

Governing equations

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$$\frac{\partial \phi}{\partial t} + \nabla \cdot (\phi \mathbf{v}_w) = \frac{r}{\rho_w}$$

- ▶ water velocity

$$\mathbf{v}_w = \mathbf{v}_i$$

- ▶ extension/
/compaction
- ▶ thermal &
phase-change
buoyancy
- ▶ bulk & shear viscous
deformation
- ▶ heat consumption/
/release
- ▶ melting/
/freezing
- ▶ **water advected by
ice**

Governing equations

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$$\mathbf{v}_w = \mathbf{v}_i - \frac{k(\phi)}{\mu_w \phi} (1-\phi) \Delta \rho \mathbf{g}$$

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- ▶ heat consumption/
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- ▶ $k(\phi)=0$: water
advected by ice
- ▶ $k(\phi)>0$: water
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- ▶ $k(\phi)=0$: water
advected by ice
- ▶ $k(\phi)>0$: water
percolates through
ice
- ▶ $\Delta \rho > 0$: water
migrates upwards

Boundary conditions

- ▶ Stokes problem

- ▶ free slip (Γ_T & Γ_S): $\mathbf{v}_i \cdot \mathbf{n} = 0$

$$\left[(1-\phi)\mu_i \left(\frac{1}{\phi}(\nabla \cdot \mathbf{v}_i)\mathbf{I} + \nabla \mathbf{v}_i + (\nabla \mathbf{v}_i)^T - \frac{2}{3}(\nabla \cdot \mathbf{v}_i)\mathbf{I} \right) \cdot \mathbf{n} \right]_t = \mathbf{0}$$

- ▶ no slip (Γ_B): $\mathbf{v}_i = \mathbf{0}$

- ▶ energy balance

- ▶ fixed temperature (Γ_T): $T = T_m$

- ▶ prescribed heat flux from silicates (Γ_B): $\nabla T = \mathbf{q}_s$

- ▶ insulating (Γ_S): $\nabla T \cdot \mathbf{n} = 0$

- ▶ water transport

- ▶ free outflow ($\Gamma_T =$ interface with the ocean)

- ▶ impermeable (Γ_B & Γ_S): $\mathbf{v}_w \cdot \mathbf{n} = 0$

Model parameters

- ▶ ice permeable to water transport above a threshold porosity ϕ_c

$$k(\phi) = \begin{cases} k_0(\phi - \phi_c)^2 & \phi \geq \phi_c \\ 0 & \phi < \phi_c \end{cases}$$

- ▶ temperature and pressure dependent ice viscosity

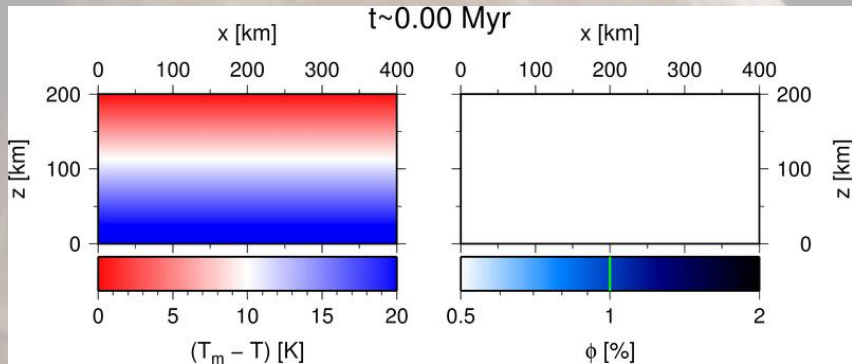
$$\mu_i = \mu_0 \exp \left[\frac{Q}{R} \left(\frac{1}{T} - \frac{1}{T_m(P)} \right) \right]$$

Numerical implementation

- ▶ 2d Finite Element code in cartesian box with $\lambda=2$
- ▶ FEniCS (<http://fenicsproject.org>), MPI parallelization
- ▶ problem split into 3 partially decoupled subproblems:
 - (1) **energy balance**: P2, SUPG regularization, Crank-Nicolson scheme
 - (2) **porosity transport**: discontinuous Galerkin, implicit Euler scheme
 - (3) **Stokes problem**: Taylor-Hood (P1-P2 for pressure-velocity)
- ▶ linear (1, 3): PETSc mumps, nonlinear (2): Newton
- ▶ mesh: 200×100 crossed diagonal triangular elements, refined at bottom and top, $\sim 10^5$ elements (best resolution less than 1km)
- ▶ adaptive time-stepping, two problems with different timescales:
 - convection: $\Delta t_c \sim \mathbf{v}_i^{-1}$
 - percolation: $\Delta t_p \sim \mathbf{v}_w^{-1}$
- ▶ computed times: $\sim 10\text{--}40$ Myr, $\Delta t \sim 10$ y, $\sim 10^6$ time steps
- ▶ typical calculation on Salomon cluster: 24 cores, 400-600 hours (depending on model parameters)
- ▶ we acknowledge support by IT4Innovations projects OPEN-10-1, OPEN-11-21

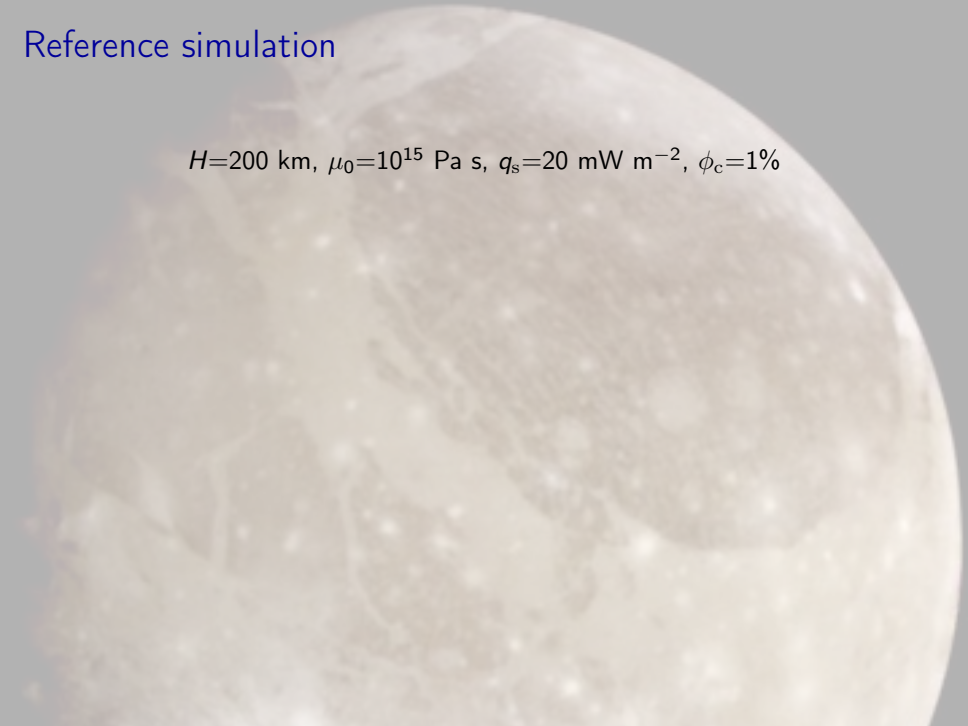
Reference simulation

$$H=200 \text{ km}, \mu_0=10^{15} \text{ Pa s}, q_s=20 \text{ mW m}^{-2}, \phi_c=1\%$$

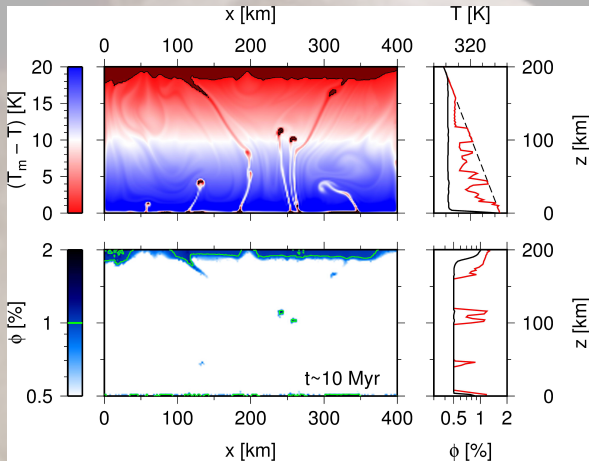


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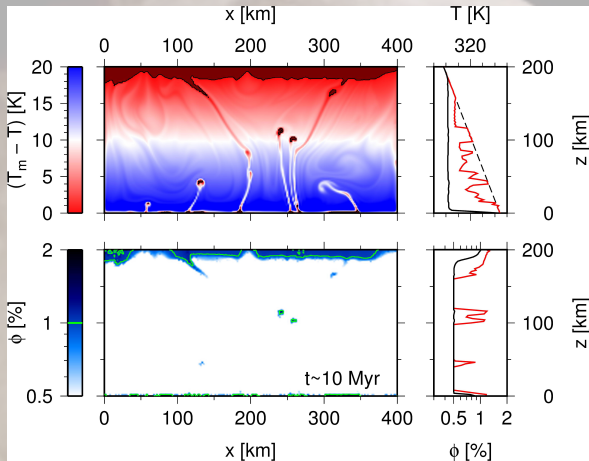


Reference simulation - statistical steady-state



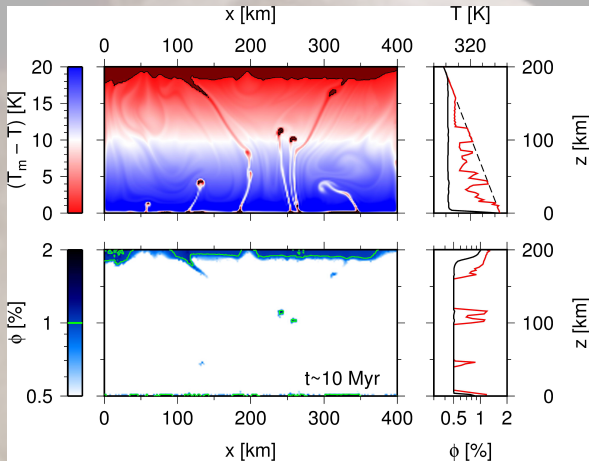
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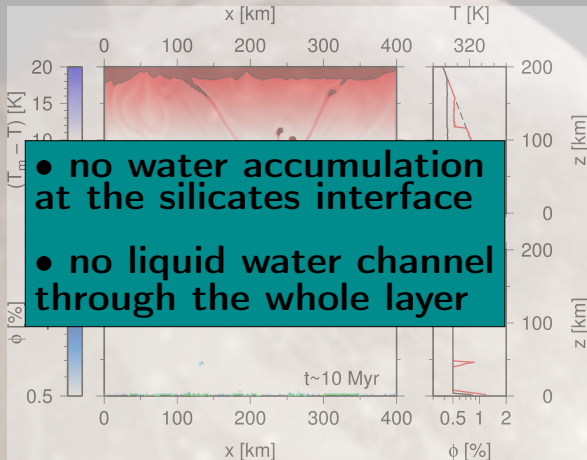
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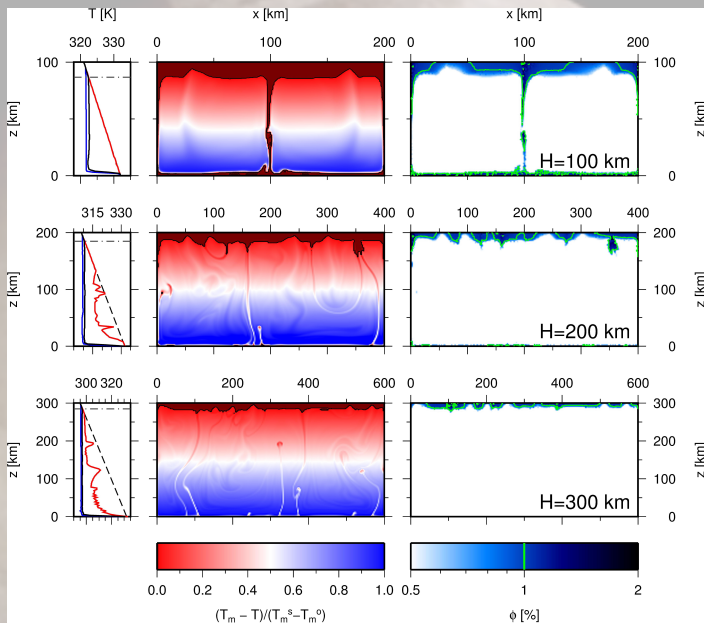
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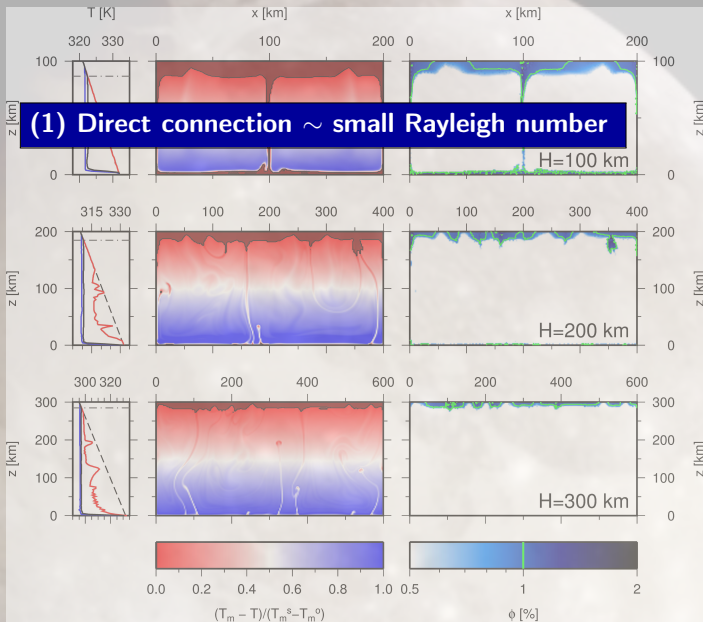


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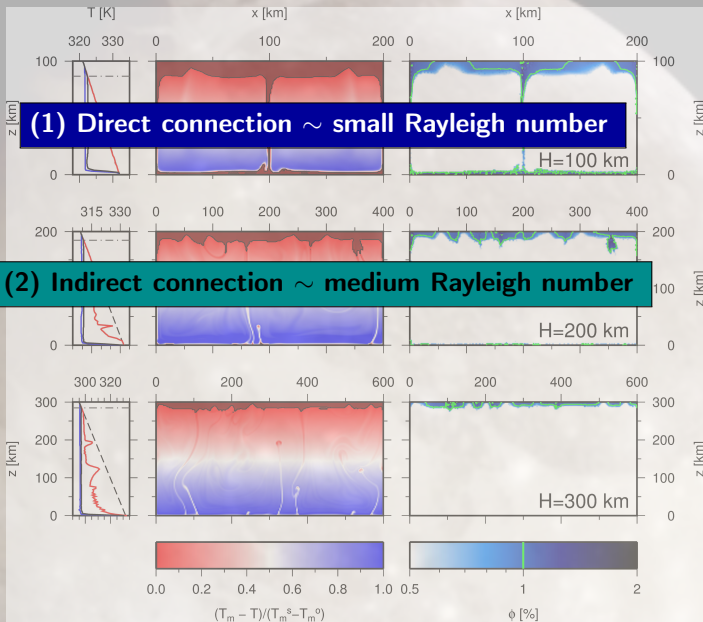
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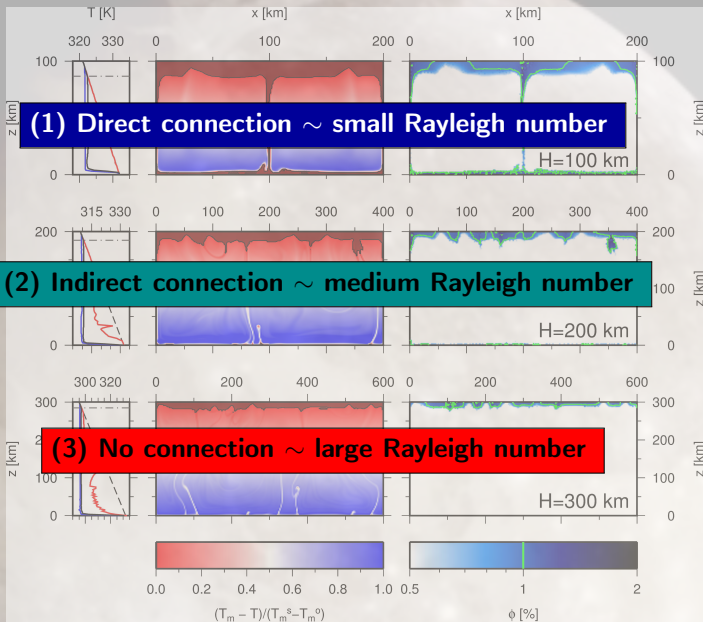
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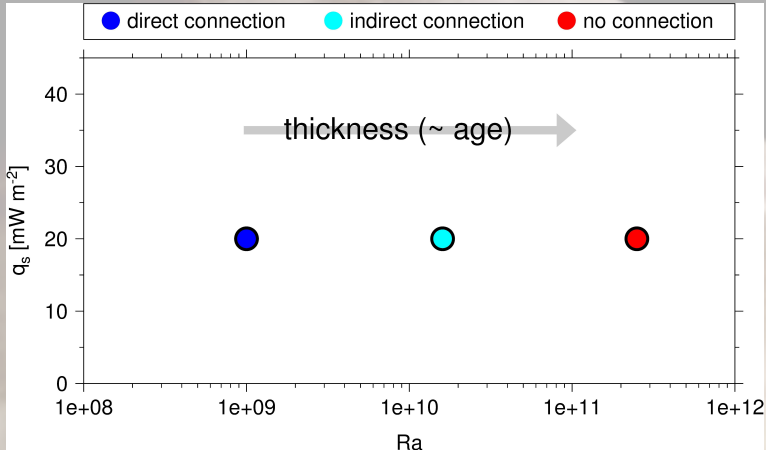
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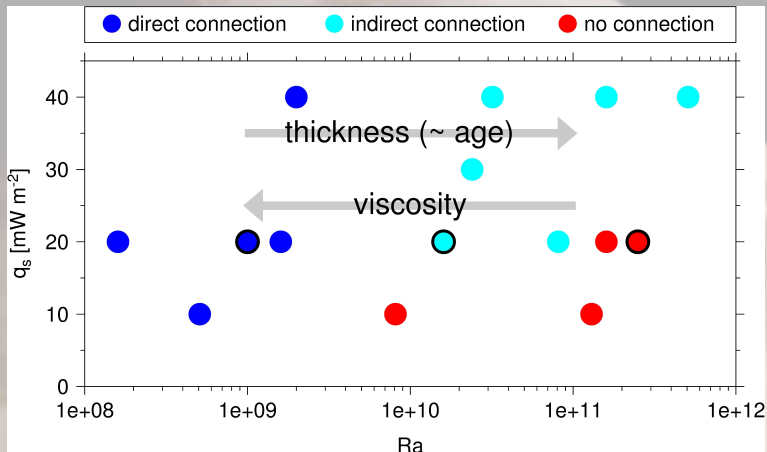


Ganymede's global thermal evolution



- ▶ HP ice layer became less permeable as Ganymede cooled down

Ganymede's global thermal evolution



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- ▶ larger viscosity \rightarrow connection even later in Ganymede's evolution
- ▶ larger silicates heat flux \rightarrow more bottom melting

Conclusions

- ▶ melting at the interface with silicates → volatiles leaching
- ▶ melt & volatiles transfer by upwelling plumes → extraction into ocean
- ▶ exchange of water and volatiles between silicate interior and ocean more probable during the early stages of Ganymede's evolution

Kalousova et al. (2018), Two-phase convection in Ganymede's high-pressure ice layer - Implications for its geological evolution, *Icarus*, **299**, 133–147.

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juice

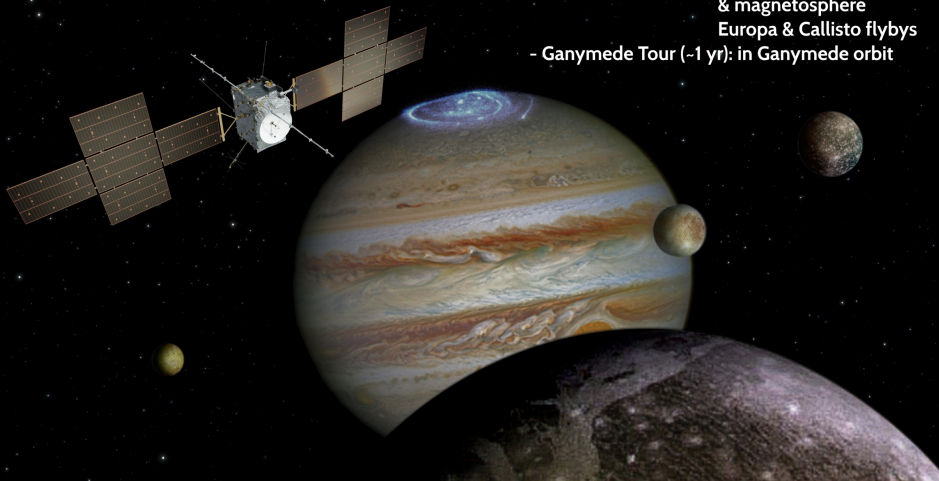


→ JUPITER ICY MOONS EXPLORER

Exploring the emergence of habitable worlds around gas giants

Two mission phases:

- Jupiter Tour (~2.5 yr): Jovian atmosphere & magnetosphere
Europa & Callisto flybys
- Ganymede Tour (~1 yr): in Ganymede orbit



Conclusions & perspectives

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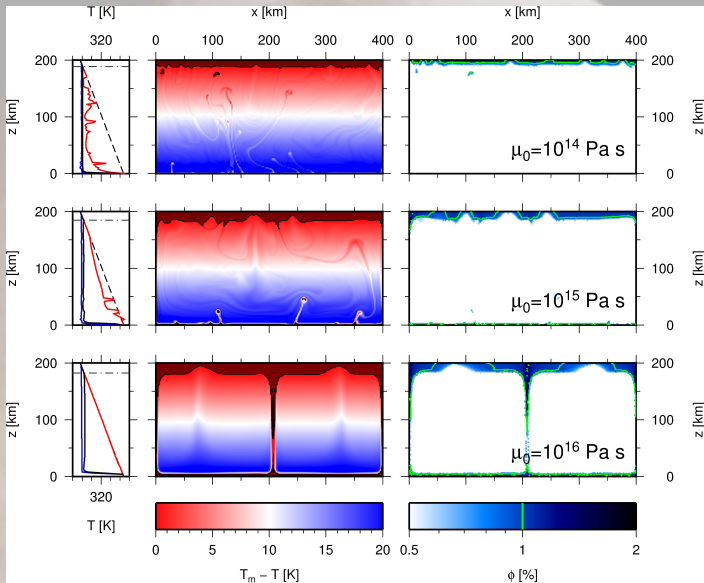
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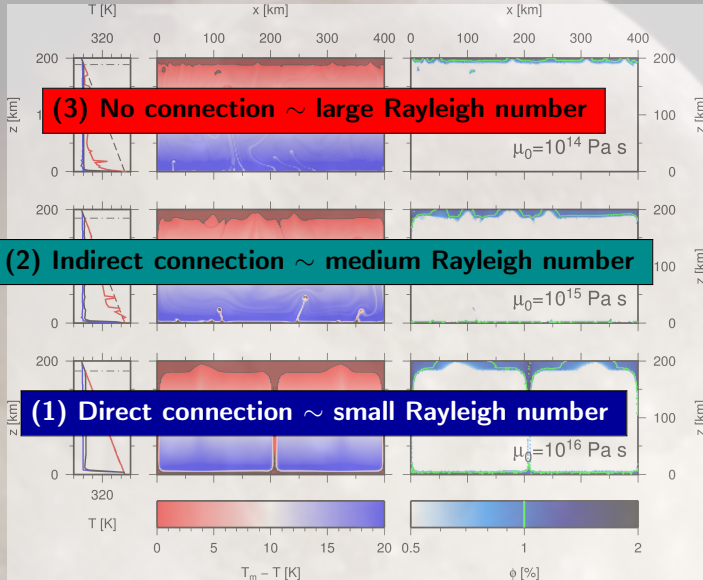
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Thank you for your
attention!

Increasing melting point viscosity μ_0 (\sim decreasing Ra)



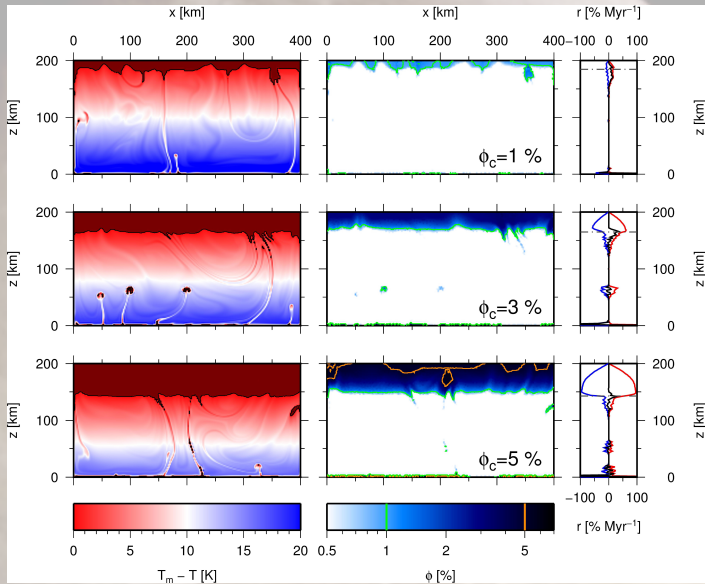
Increasing melting point viscosity μ_0 (\sim decreasing Ra)



Increasing percolation threshold ϕ_c (top to bottom)

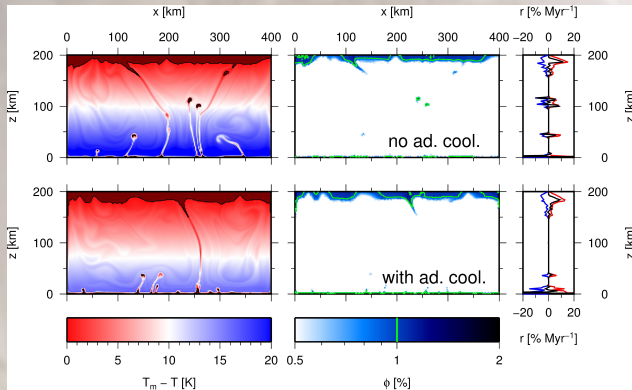
→ thicker top temperate layer

→ more melt at the interface with ocean



Effects of adiabatic heating/cooling

- ▶ adiabatic gradient: $\left(\frac{\partial T}{\partial r}\right)_s = -\frac{\alpha T g}{c_p} \sim -0.026 \text{ K km}^{-1}$
- ▶ temperature gradient across the layer: $\left(\frac{\partial T}{\partial r}\right) \sim -0.115 \text{ K km}^{-1}$



- ▶ cooling due to decompression \rightarrow less melting in the upwelling plumes
 - ▶ warming up of descending material due to compression
- \rightarrow warmer interior \sim enhanced potential for melting at the silicates interface