

Intrinsic Friction

A. Cammarata

TMDs

TMDs as Dry Lubricants

Simulations of MX<sub>2</sub> TMDs

Classical Simulations

Intrinsic Friction

Quantum Simulations

Dynamic Analysis

Cophonicity

Ti:MoS<sub>2</sub>

NMTA

TMDs Under Load

Finite *n*-Layer Systems

Charged Systems

Conclusions and Future Directions

# Harness Nanoscale Friction in Transition Metal Dichalcogenides

## 1st Users' Conference

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Czech Technical University in Prague - Advanced Materials Group  
Support: Czech Science Foundation, projects 16-11516Y and 17-24164Y  
IT4Innovations Centre of Excellence

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# Advanced Materials Group



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# Transition Metal Dichalcogenides

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 $\text{Ti}:\text{MoS}_2$ 

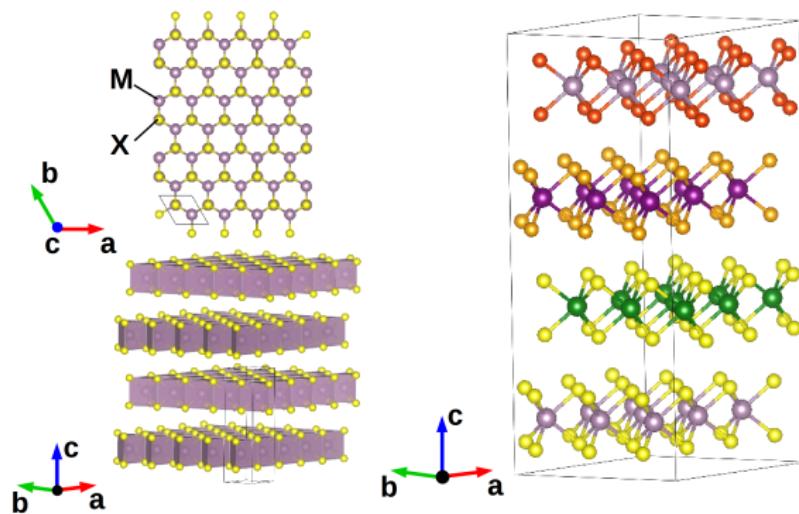
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$\text{MX}_2$  Transition Metal Dichalcogenides: flexible chemistry ( $\text{M} = \text{Mo}, \text{W}, \text{Ti}, \text{Zr}, \text{V} \dots \text{X} = \text{S}, \text{Se}, \text{Te}$ )  $\Rightarrow$  easy property tuning and wide applicability (lubrication, catalysis, energy storage, electronic devices)

[*Nature Chemistry* **5**, 263 (2013); *Nature* **499**, 419 (2013)]

# Transition Metal Dichalcogenides as Dry Lubricants

Dry lubricants: mandatory in extreme conditions like high temperature or under vacuum

- ▶ Transition metal dichalcogenides (TMDs)

Experimental discovery of new TMD-based materials with optimal characteristics is challenging and time-consuming, requiring development cycles that include:

- ▶ candidate material identification
- ▶ testing
- ▶ structural optimization

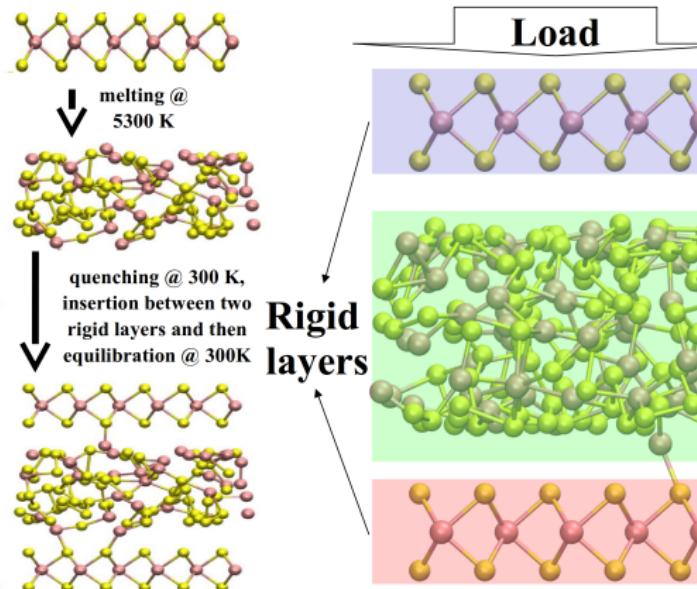
Computational techniques expedite the process by narrowing the composition-structure phase space for experimental exploration to only the most favorable compounds:

⇒ efficient design of novel TMD materials with improved or new frictional properties

# Classical simulations

[P. Nicolini, T. Polcar et al. submitted to *Nanoscale*]

## Crystal formation from amorphous phase



- Time step 0.1 fs
- 9.3 x 10.8 nm
- 10404 atoms
- 2-4 ns MD run
- T 3000-6000 K
- Load: 50 GPa
- 6 coating densities
- AIREBO potential
- Speed 0, 1, 5 m/s

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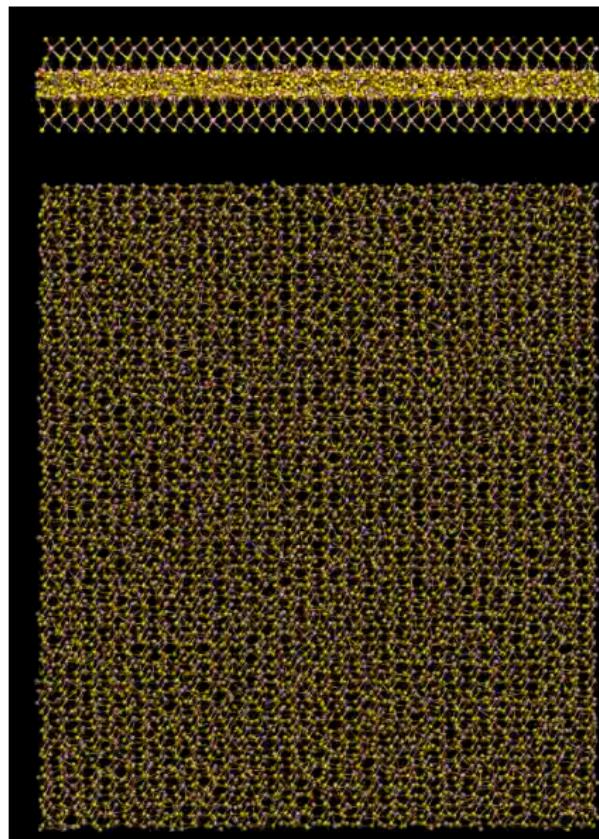
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# Classical simulations

[P. Nicolini, T. Polcar et al. submitted to *Nanoscale*]



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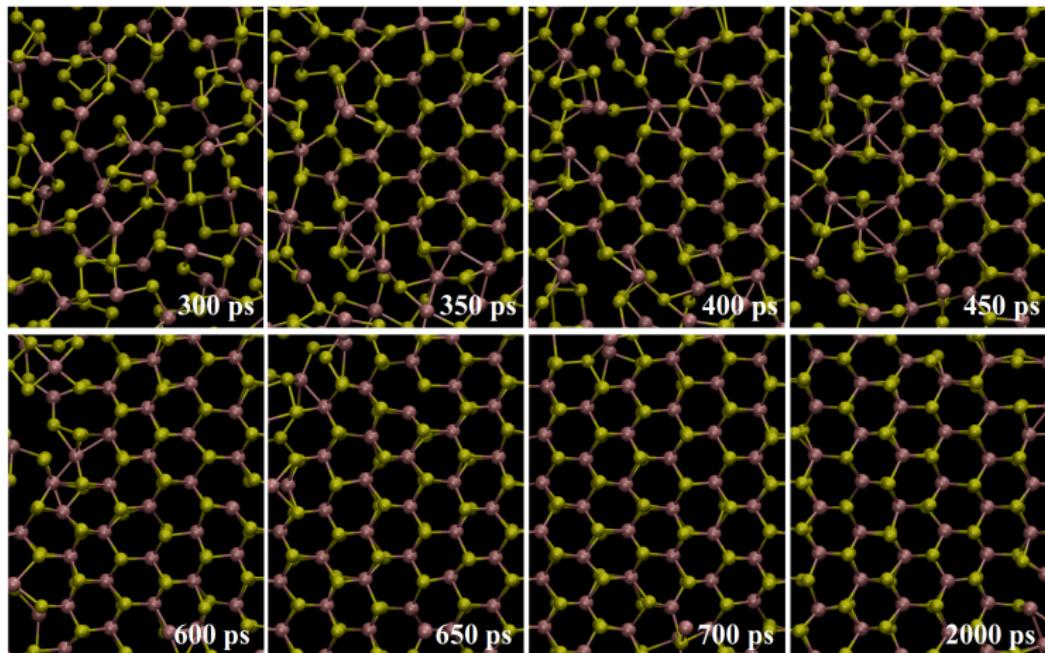
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# Classical simulations

[P. Nicolini, T. Polcar et al. submitted to *Nanoscale*]



# Intrinsic Friction

**Friction:** **macroscopic** effect of a resultant force that resists sliding or rolling of one object over another

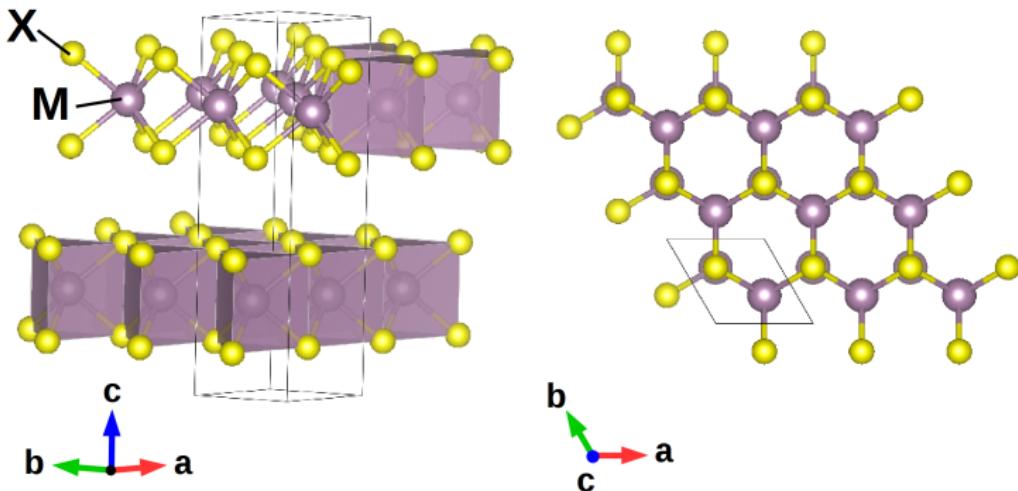
**Microscopic Friction:** friction generated at the **atomic scale** by the relative motion of few adjacent atom layers in the presence of structural irregularities (dislocations, layer truncations, grain boundary orientation...)

**Intrinsic Friction:** **microscopic friction** due to only local electronic (atomic type) and structural features (geometry) **without irregularities**  $\Rightarrow$  a property *intrinsic* of the system

## Quantum Mechanical Modeling of Intrinsic Friction in TMDs

# Quantum Mechanical Models of MX<sub>2</sub> TMDs

MX<sub>2</sub> — M = Mo, W; X = S, Se, Te (*P6<sub>3</sub>/mmc*)



Ab Initio simulations allow to identify and disentangle the electro-structural features that determine the intra- and inter-layer motions affecting the intrinsic friction.

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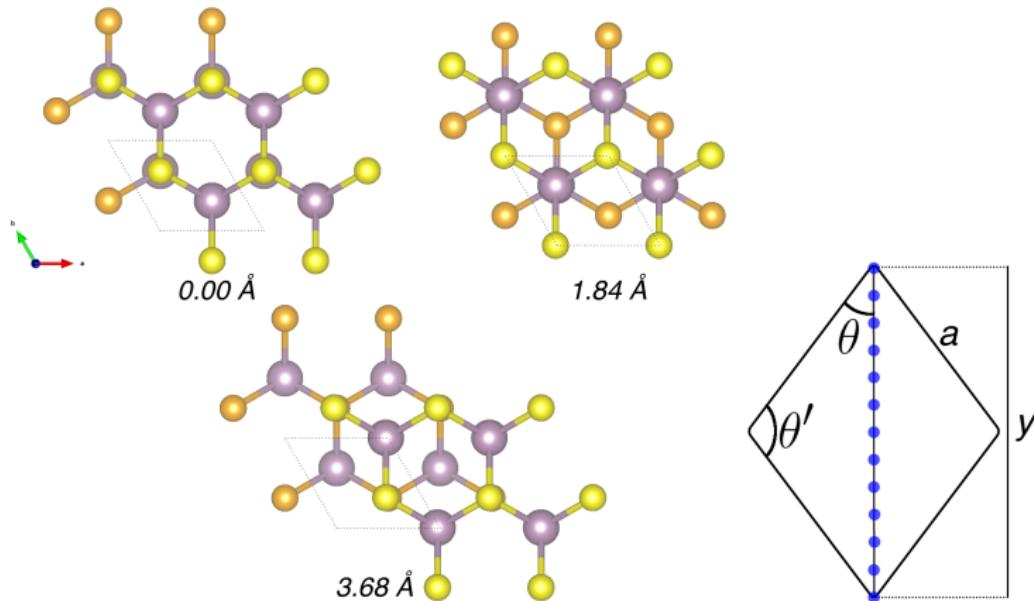
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# Quantum Mechanical Models of MX<sub>2</sub> TMDs

[B. J. Irving, P. Nicolini, T. Polcar, *Nanoscale* 9, 5597 (2017)]

Most important geometric configurations for MX<sub>2</sub> bilayers:



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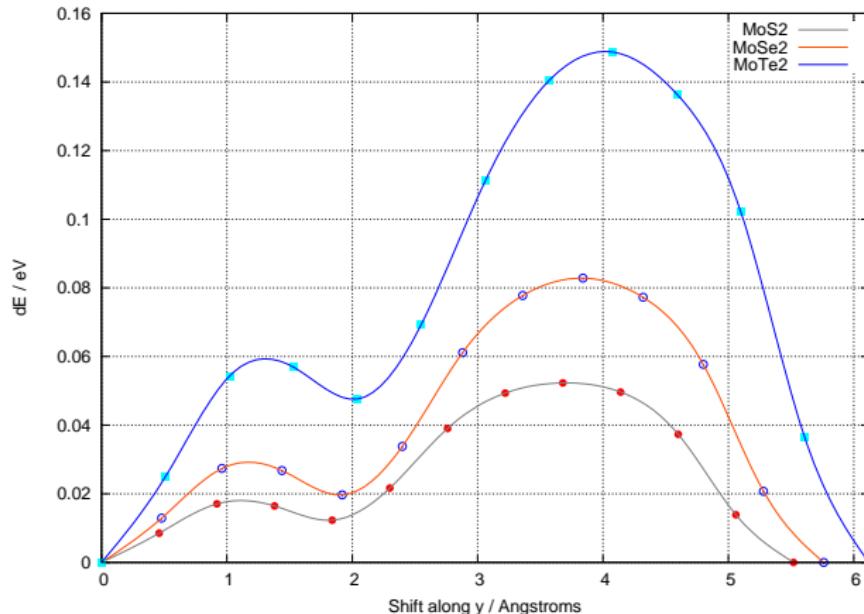
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# Quantum Mechanical Models of MX<sub>2</sub> TMDs

[B. J. Irving, P. Nicolini, T. Polcar, *Nanoscale* 9, 5597 (2017)]

Potential energy profiles for MoX<sub>2</sub> bilayers:



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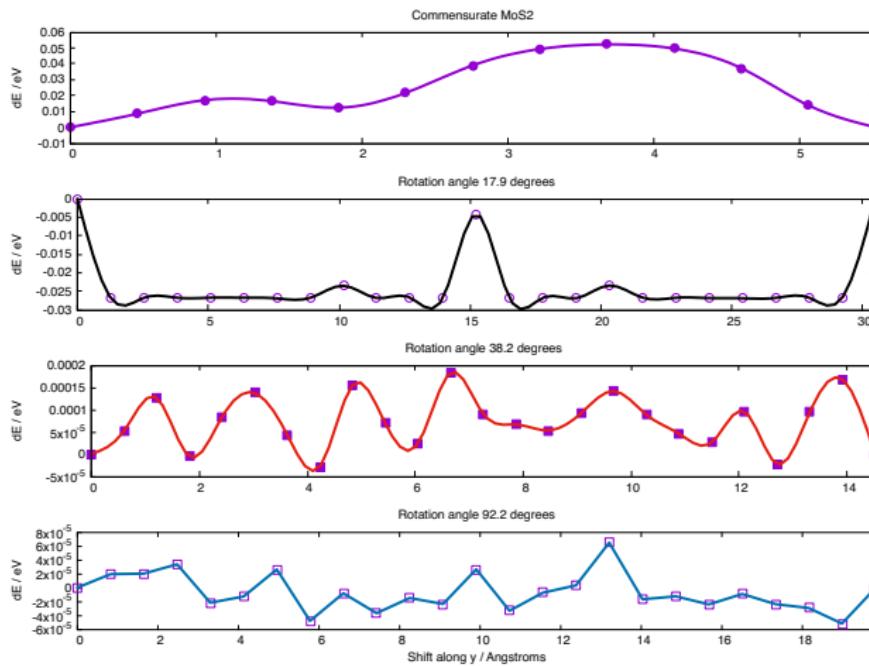
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# Quantum Mechanical Models of MX<sub>2</sub> TMDs

[B. J. Irving, P. Nicolini, T. Polcar, *Nanoscale* 9, 5597 (2017)]

Incommensuration effects: a greater degree of incommensuration results in a significant easier inter-plane shear



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# Disentangling Electro-Structural Features

In terms of the classical picture, atoms are called back by a harmonic restoring force with associated frequency  $\omega = \sqrt{k_f/\mu}$

lower frequency  $\Rightarrow$  sliding promoted

Goal: to improve intrinsic friction, that is, to lower the sliding-related frequencies

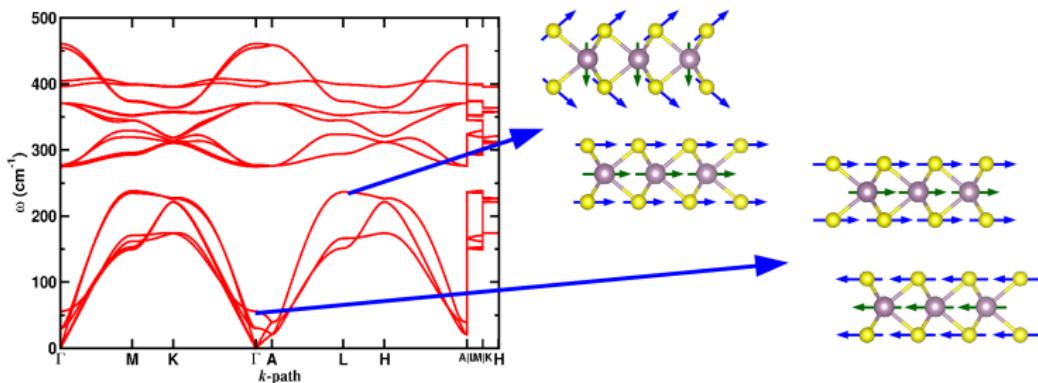
We parameterize the electronic and structural features by means of simple descriptors:

- ▶ Dynamic features: Cophonicity
- ▶ Electronic density: Covalency, Orbital Polarization
- ▶ Structural Distortions: Group Theory

The screenshot shows a Wikipedia page for "Cophonicity". The page header includes the Wikipedia logo, a globe icon, and navigation links for "Article", "Talk", "Read", "Edit", "View history", and "Search". The main title "Cophonicity" is bolded, followed by the text "From Wikipedia, the free encyclopedia". The page content explains that cophonicity is a measure of the overlap of atomic contributions to vibrational frequencies, specifically for the A-B atomic pair. It is described as a metric used to parametrize dynamical interactions.

# Dynamic Analysis

## Phonon Band Structure



- ▶ Identification of the vibrational modes related to layer sliding
- ▶ Each mode corresponds to a harmonic restoring force with force constant  $k_f$ :  $\omega = \sqrt{k_f/\mu}$
- ▶ The lower the frequency, the weaker the restoring force, the higher the atomic displacement amplitude

Optimal  $\omega \Rightarrow$  optimal layer sliding

We need a descriptor (metric) to quantify how changing the atomic type affects specific phonon frequencies

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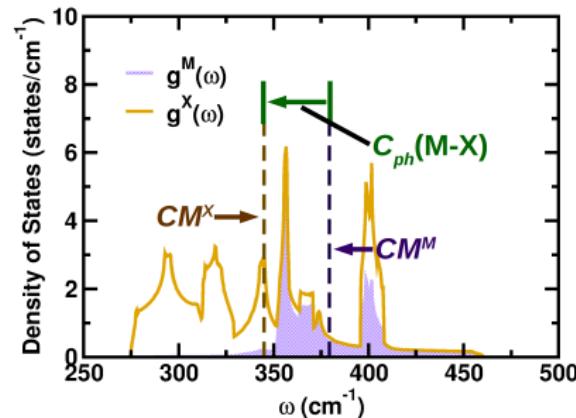
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# The Cophonicity Metric

[A. Cammarata, T. Polcar, *Inorg. Chem.* **54**, 5739 (2015)]

The phonon density of states can be decomposed into atomic contributions of the M-X atomic pair, which is the smallest unit that generates the dynamical interactions



The relative position  $C_{ph}(M-X)$  of  $CM^M$  with respect to  $CM^X$  is  
$$C_{ph}(M-X) = CM^M - CM^X$$

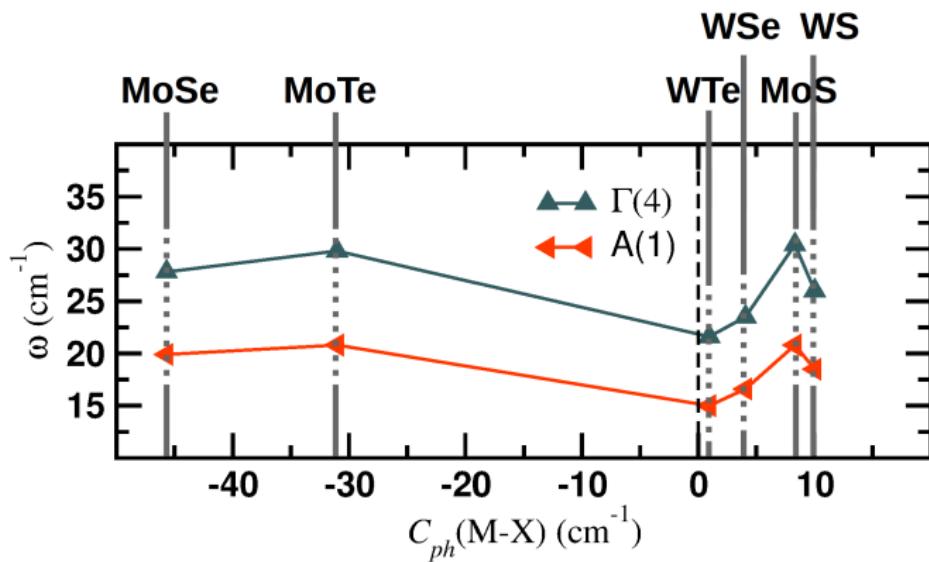
We call  $C_{ph}(M-X)$  the **cophonicity** of the M-X atomic pair

Cophonicity is a measure of the overlap of the single atomic contributions to a specific frequency range.

# Results

## Tuning Frequency with Cophonicity

Mode frequency is minimum when  $C_{ph} \simeq 0 \text{ cm}^{-1}$



How can we control M-X pair cophonicity?

# Results

## Tuning Cophonicity with Covalency and Structural Distortions

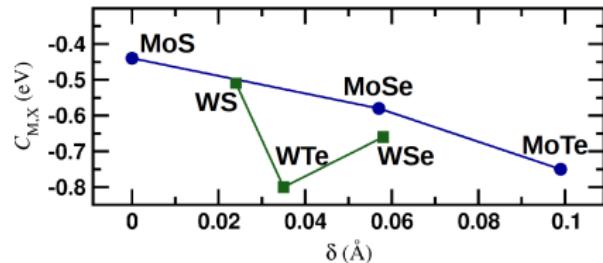
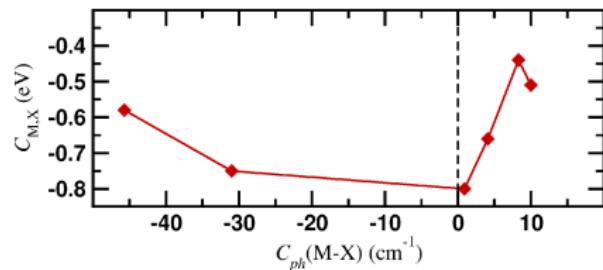
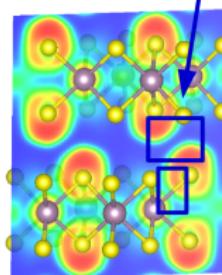
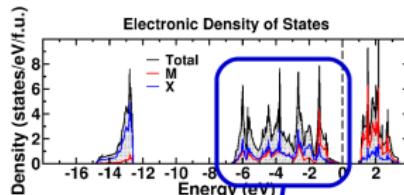
Cophonicity is related to M–X bond covalency and the structural distortions

### M–X bond covalency $C_{M,X}$

A. Cammarata et al.

J. Chem. Phys. **141**, 114704 (2014)

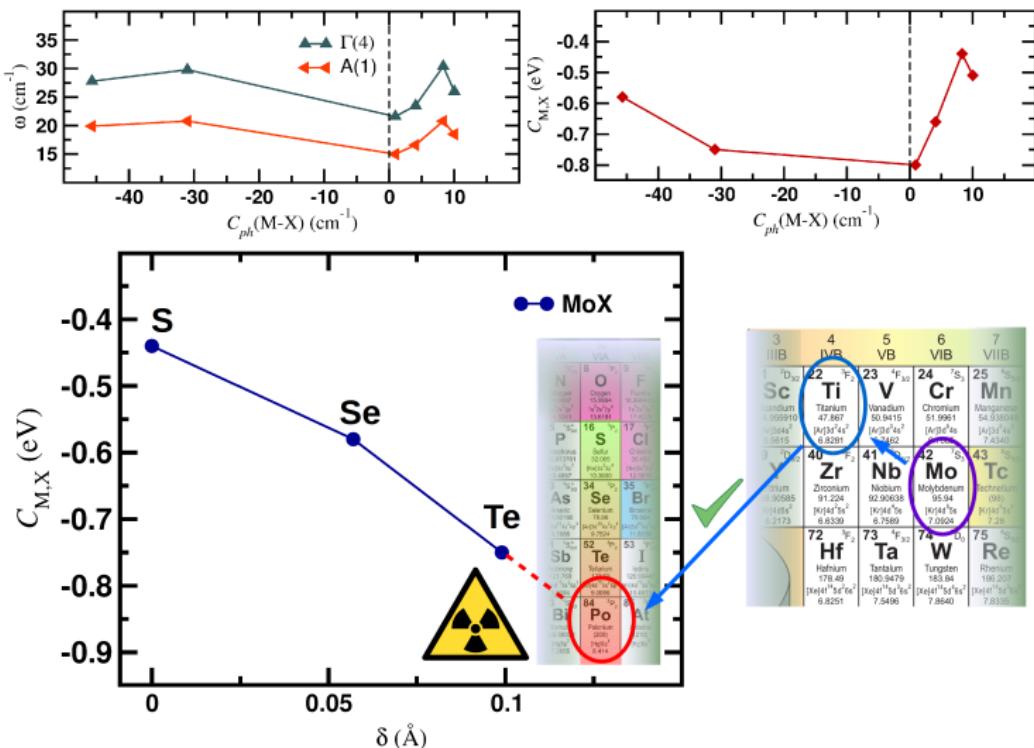
Chem. Mater. **26**, 5773 (2014)



Let's now design a new TMD!

# Designing a new TMD

[A. Cammarata, T. Polcar, *Inorg. Chem.* **54**, 5739 (2015)]



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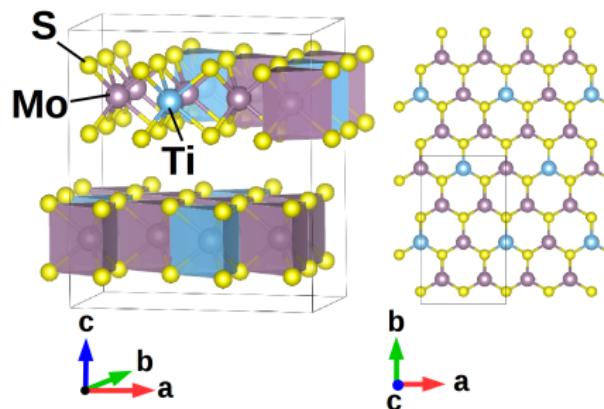
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# The Ti:MoS<sub>2</sub> System

[A. Cammarata, T. Polcar, *Inorg. Chem.* **54**, 5739 (2015)]



- The **most stable** among the explored MX<sub>2</sub>, comparable to MoS<sub>2</sub>
- **Lowest sliding frequencies** expected, comparable to WSe<sub>2</sub> and WTe<sub>2</sub>

System	$\Delta H_f$ (eV/f.u.)	$\delta$ (Å)	$C_{M,X}$ (eV)	$C_{ph}(M-X)$ (cm <sup>-1</sup> )	$\omega(\Gamma)_{4-5}$ (cm <sup>-1</sup> )	$\omega(A)_{1-4}$ (cm <sup>-1</sup> )
MoS	0.00	0.00	-0.44	8.3	30	21
MoSe	2.79	0.06	-0.58	-45.7	28	20
MoTe	5.94	0.10	-0.75	-31.0	30	21
WS	1.23	0.02	-0.51	10.0	26	18
WSe	4.27	0.06	-0.66	4.1	24	17
WTe	7.63	0.03	-0.80	0.9	22	15
Ti:MoS	0.04	0.17	-0.76	4.6	$\omega(\Gamma)_4$ 22 $\omega(\Gamma)_5$ 23	$\omega(A)_{1-2}$ 15 $\omega(A)_{3-4}$ 16

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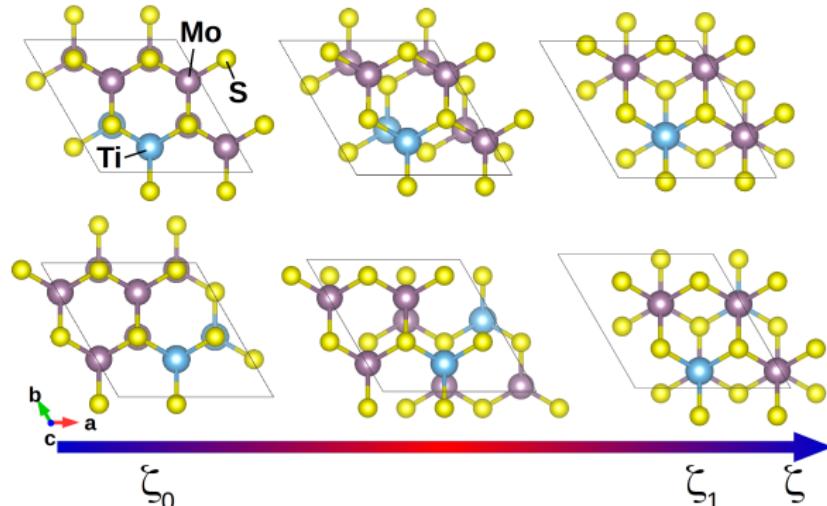
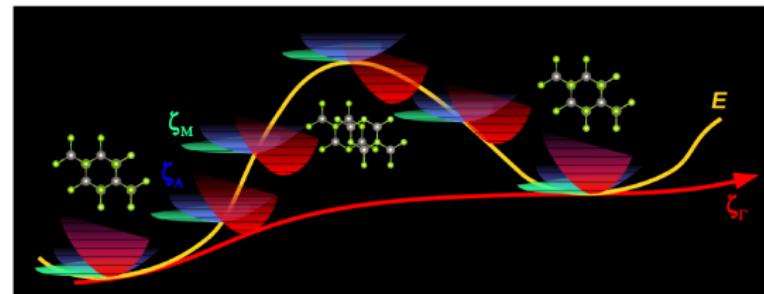
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# Normal-Modes Transition Approximation

[A. Cammarata, T. Polcar, *Phys. Rev. B* **96**, 085406 (2017)]



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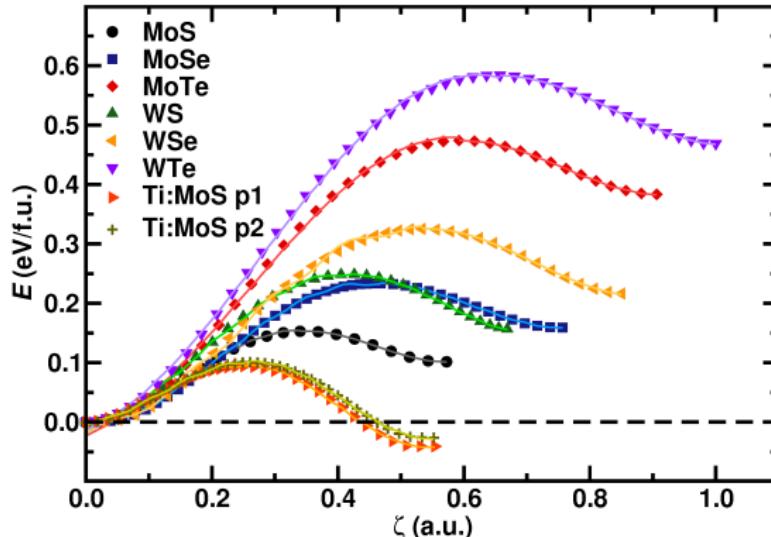
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# Normal-Modes Transition Approximation

[A. Cammarata, T. Polcar, *Phys. Rev. B* **96**, 085406 (2017)]

Along a general configurational coordinate  $\zeta$ , the total potential energy  $E$  can be approximated as a sum of contributions from all the  $K$  normal modes:

$$E(\zeta) = \sum_{k=1}^K \frac{1}{4} \omega_k^2 C_k^2 \left[ 1 - \cos \left( \frac{\pi}{C_k} \zeta - \theta_k \right) \right] \quad (1)$$



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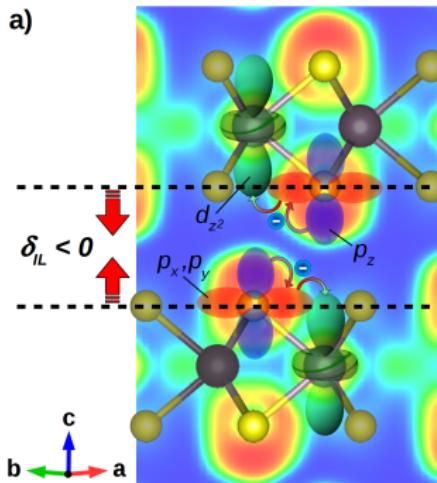
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# MX<sub>2</sub> Under Load

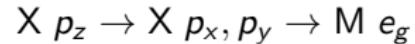
[A. Cammarata, T. Polcar, RSC Adv. 5, 106809 (2015)]

a)



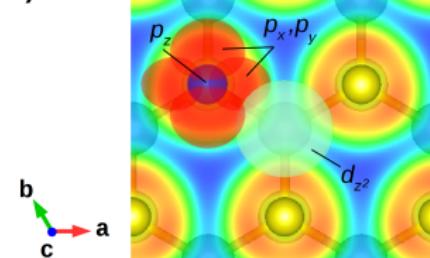
▶ Interlayer charge distribution is determined by the  $p_z$  orbital of the X atoms

▶ Under load, electronic charge flows from inter- to intralayer region via  $p_x$  and  $p_y$  orbitals of X atom:



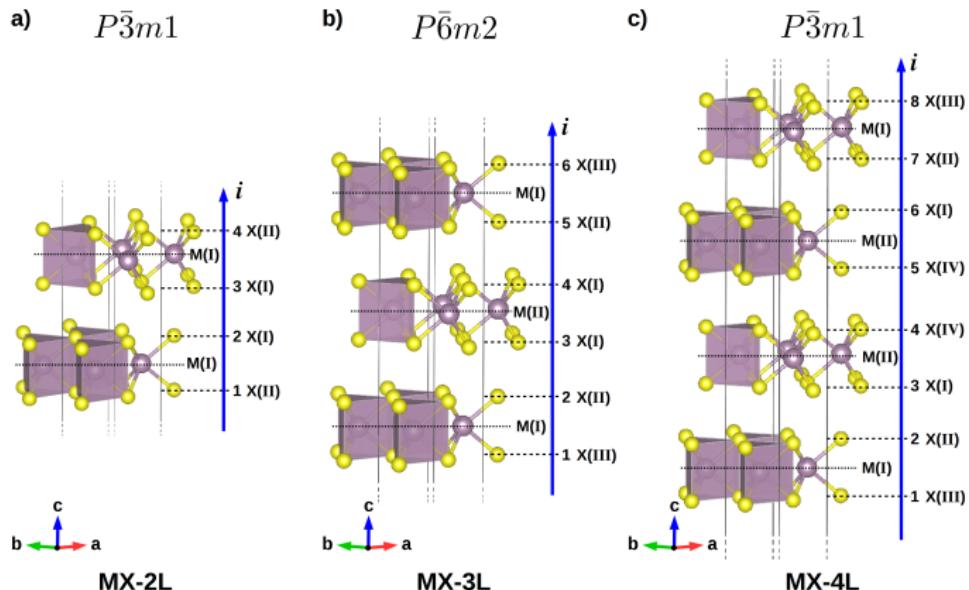
▶ Cophonicity is independent of the load and restricts the frequency variation against the applied load

b)



# MX<sub>2</sub> *n*-layer systems

[A. Cammarata, T. Polcar, *Phys. Chem. Chem. Phys.* **18**, 4807 (2016)]



Cophonicity  $\Rightarrow$  Inter-layer orb. pol.  $\Rightarrow$  Electro-vibrational coupling

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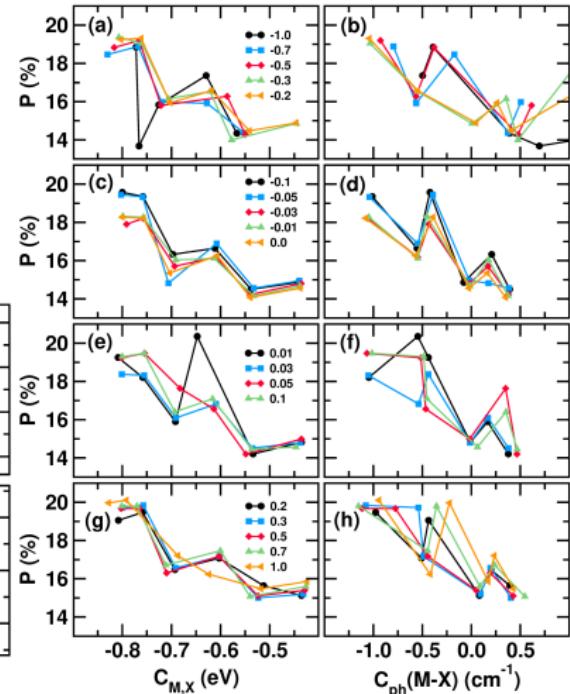
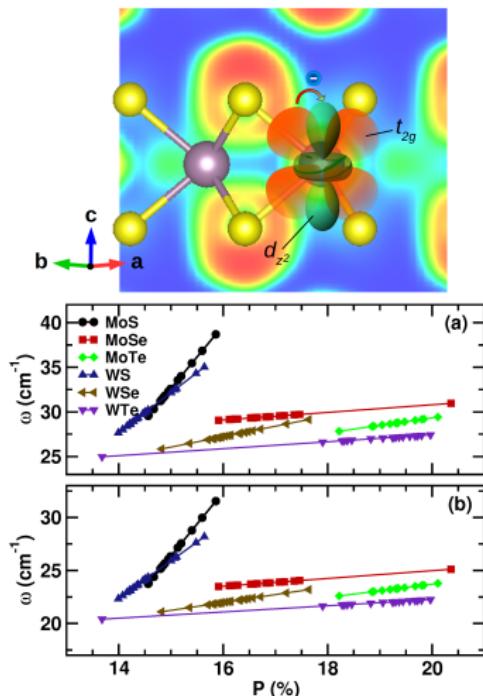
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# Charged MX<sub>2</sub> Systems

[A. Cammarata, T. Polcar, *Nanoscale* 9, 11488 (2017)]



Cophonicity and covalency as a knob to control electron-phonon coupling

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# Conclusions and Future Directions

## Results:

- Crystal MoS<sub>2</sub> formation from amorphous phase by temperature quenching
- The greater the incommensuration, the easier the inter-plane shear
- New lattice dynamic metric proposed: *cophonicity*
- New promising tribological material proposed: Ti:MoS<sub>2</sub>
- ◆ NMTA method decomposes MEPs into phonon contributions
- ◆ External Load and perturbation of charge neutrality affect sliding-related vibrational modes
- ◆ Cophonicity and covalency are a knob to control the electron-phonon coupling

## Future Directions:

- Activation/deactivation of selected modes to favor/disfavor layer sliding, exfoliation, internal heat dissipation...
- ◆ Applications beyond tribology: Metal-Insulator transitions, linear dielectric response, nonlinear Second Harmonic Generation, phase matchability, ionic conduction...

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Thank you very much  
for your kind attention

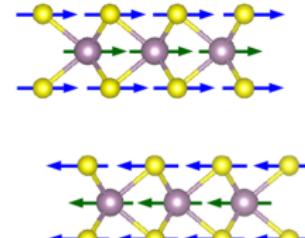
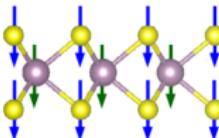
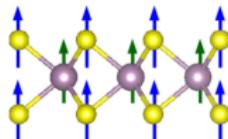
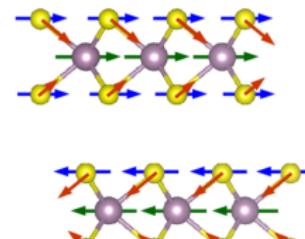
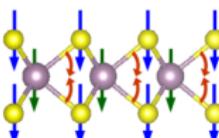
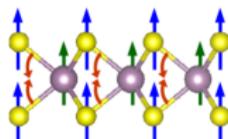






# Dynamic Analysis

## Example of sliding-related modes

**a)** $\Gamma(4-5)$ 
**b)** $\Gamma(6)$ 
**c)** $A(1-4)$ 
**d)** $A(5-6)$ 


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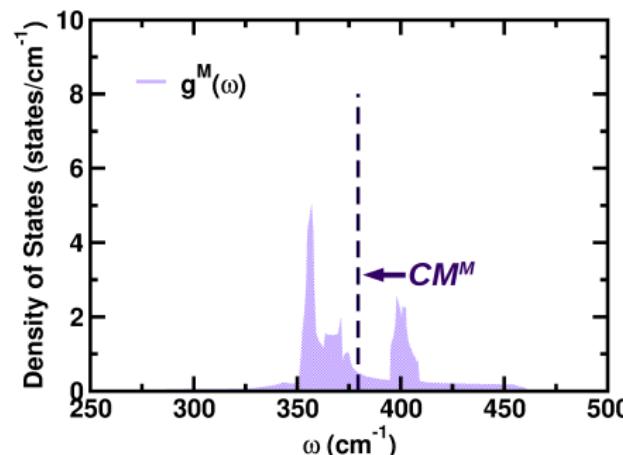
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# The Cophonicity Metric

[A. Cammarata, T. Polcar, *Inorg. Chem.* **54**, 5739 (2015)]

In a specific range  $[\omega_0, \omega_1]$ , the center mass  $CM^M$  of  $g^M(\omega)$

$$\text{is defined as } CM^M = \frac{\frac{\omega_1}{\omega_0} \int_{\omega_1}^{\omega_0} \omega g^M(\omega) d\omega}{\int_{\omega_0}^{\omega_1} g^M(\omega) d\omega}$$



We call  $\int_{\omega_0}^{\omega_1} g^M(\omega) d\omega$  the **phonicity** of the M atom

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Ti:MoS<sub>2</sub>

NMTA

TMDs Under Load

Finite  $n$ -Layer Systems

Charged Systems

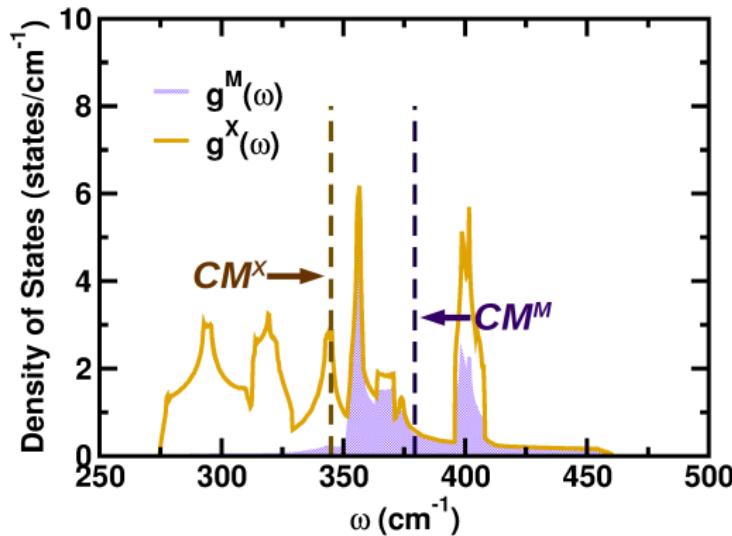
Conclusions and

Future Directions

# The Cophonicity Metric

[A. Cammarata, T. Polcar, *Inorg. Chem.* **54**, 5739 (2015)]

In general, M and X atomic contributions are centered around distinct phonon frequencies



Intrinsic Friction

A. Cammarata

TMDs

TMDs as Dry Lubricants

Simulations of MX<sub>2</sub> TMDs

Classical Simulations

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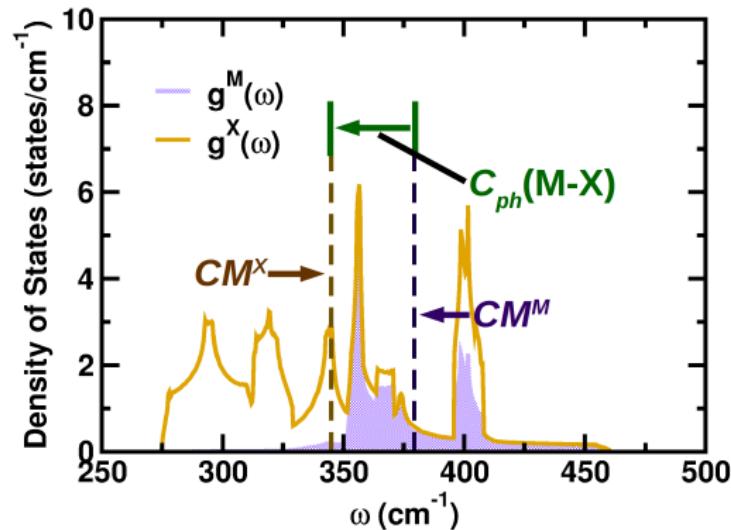
Charged Systems

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# The *Cophonicity* Metric

[A. Cammarata, T. Polcar, *Inorg. Chem.* **54**, 5739 (2015)]

The relative position  $C_{ph}(M-X)$  of  $CM^M$  with respect to  $CM^X$  is  
$$C_{ph}(M-X) = CM^M - CM^X$$



We call  $C_{ph}(M-X)$  the **cophonicity** of the M-X atomic pair

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# Computational Details

- ▶ Geometries: Hexagonal  $P6_3/mmc$  (SG# 194)
- ▶ closed shell configuration
- ▶ Density Functional Theory, General Gradient Approximation (DFT-GGA): Perdew-Burke-Ernzerhof (PBE) (Phys. Rev. Lett. **77**, 3865 (1996))
- ▶ DFT-D2 Grimme correction (J. Comp. Chem. **27**, 1787 (2006))
- ▶ Vienna Ab-initio Simulation Package (VASP)
- ▶ Plane-wave cut-off: 700 eV
- ▶ PHONOPY package
- ▶ Monkhorst-Pack mesh:  $7 \times 7 \times 5$  for VASP,  $41 \times 41 \times 41$  for PHONOPY

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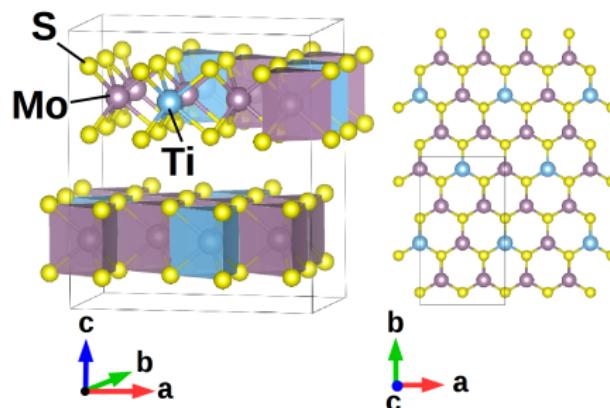
Conclusions and

Future Directions

# The Ti:MoS<sub>2</sub> System

Cmcm (SG 63), relaxed structure

[A. Cammarata, T. Polcar, *Inorg. Chem.* **54**, 5739 (2015).]



Lattice parameters:  $a = 6.46321$ ,  $b = 11.19461$ ,  $c = 12.56747$

Mo	4c	0	0.16767	$\frac{1}{4}$
Ti	4c	0	0.41744	$\frac{1}{4}$
Mo	8g	$\frac{1}{4}$	0.16767	$\frac{1}{4}$
S	8f	0	0.58487	0.12745
S	8f	0	0.08424	0.12685
S	16h	0.74866	0.33385	0.12831