

Ultrafast demagnetization dynamics observed at $M_{2,3}$ edges of hcp Co - first principles calculations

D. Legut¹, K. Carva², P.M. Oppeneer³
in collaboration with the group of Prof. M. Murnane⁴
(D. Zusin⁴, E. Turgut⁴, P. Grychtol⁴, etc.)



IT4Innovations
national supercomputing
center



- ¹IT4Innovations, VSB-Technical University of Ostrava, Ostrava, Czech Republic
²Department of Condensed Matter Physics, Charles University, Prague, Czech Republic
³Department of Physics and Astronomy, Uppsala University, Uppsala, Sweden
⁴Department of Physics and JILA, University of Colorado, Boulder, Colorado, USA

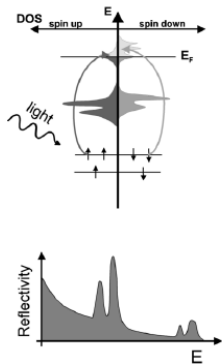
Motivation: magneto-transport linear and quadratic in \vec{M}

	conductivity	magneto-optics	X-ray
$\varepsilon_{ij}^{(1)}, i \neq j$	Hall effect	MOKE, magnetic circular dichroism and birefringence	XMCD
$\varepsilon_{ij}^{(2)}, i \neq j$	longitudinal Hall (quadratic Hall)	quadratic MOKE	quadratic XMCD
difference of diagonals, $\varepsilon_{ii}^{(2)} - \varepsilon_{jj}^{(2)} (i \neq j)$	AMR	Voigt effect, magnetic linear dichroism (MLD), Schäfer-Hubert	XMLD

- different photon energy ranges: dc (i.e. $\omega = 0$), visible, X-ray
- what part of permittivity/conductivity tensor is measured (diag./off-diag.)
- but: symmetry arguments equal for **all** quadratic magneto-transport effects
- presented arguments are equally valid both for permittivity ε_{ij} and conductivity σ_{ij} tensors as they are related by $\sigma_{ij} = \delta_{ij} + (i/\omega)\varepsilon_{ij}$.

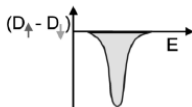
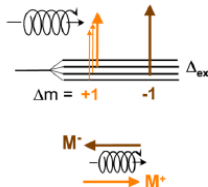
Electronic structure probed by magneto-optical spect.

general spectroscopy



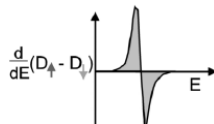
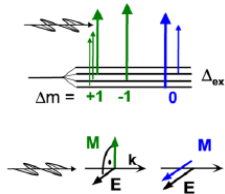
XMCD

circularly polarized light



XMLD

linearly polarized light

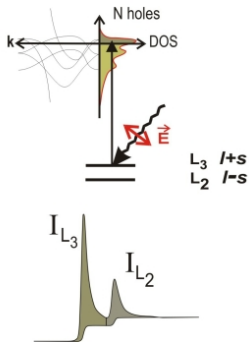


➔ most sensitive to magneto-crystalline anisotropy: XMLD

JMMM 272, 2146 (2004), PRL 104, 187401 (2010), PRB 84, 054448 (2011)

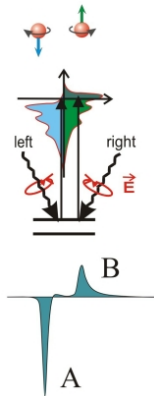
Spin and Orbital Moments: X-Ray Magnetic Circular Dichroism

(a) d-Orbital Occupation



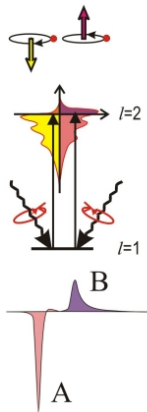
$$N_h = \langle I_{L_3} + I_{L_2} \rangle / C$$

(b) Spin Moment



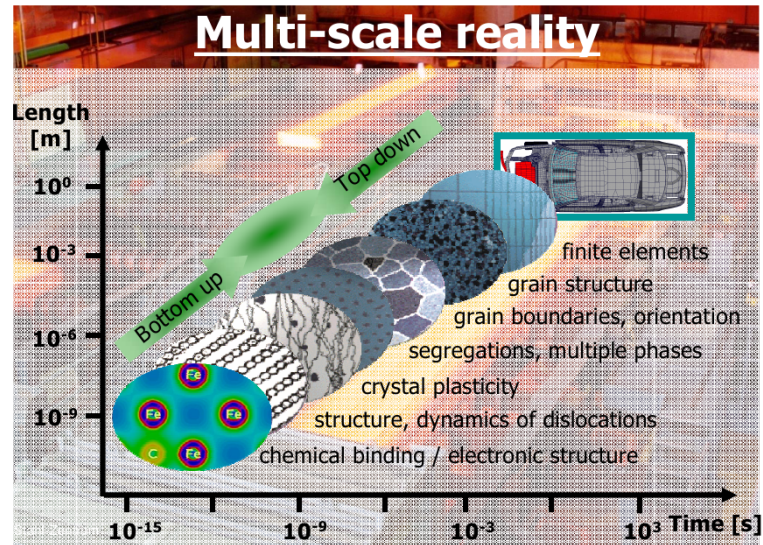
$$m_s = \mu_B \langle -A + 2B \rangle / C$$

(c) Orbital Moment



$$m_o = \mu_B \langle A + B \rangle / 3C$$

<http://ssrl.slac.stanford.edu/stohr/xmcd.htm>



First-principles calculations - utilizing HPC

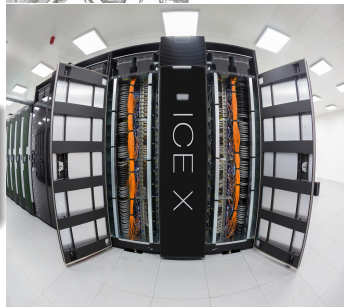
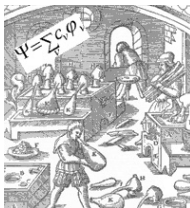
Many chemical and physical properties could be calculated from first-principles (ab initio) electronic structure calculations based on solving of:

**Fundamental quantum theory,
Schrödinger/Dirac equations
&
Density functional theory***

The only input:

- atomic number of constituents
- some structural information

* in 1998 awarded by Nobel prize in Chemistry, W. Kohn & J. Pople



Quantum-mechanical description

Many-body ($\sim 10^{23}$ per cm^3) Schrödinger equation:

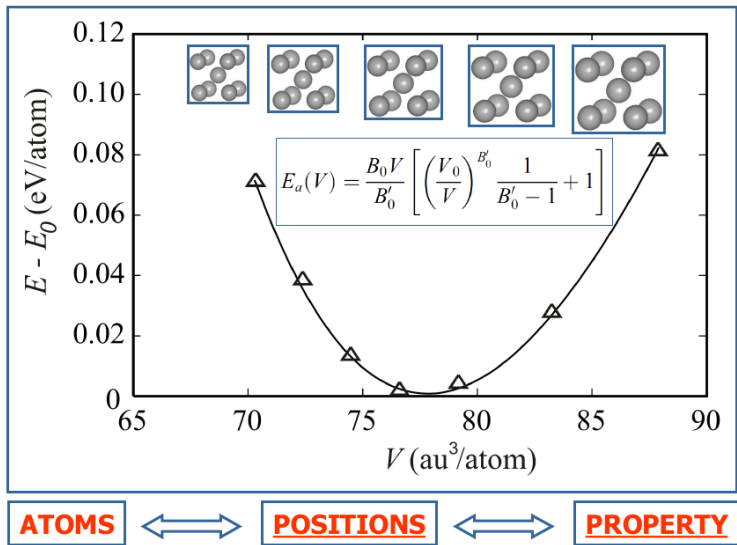
$$H \Psi(\{\mathbf{r}_i\}, \{\sigma_i\}) = E \Psi(\{\mathbf{r}_i\}, \{\sigma_i\})$$



$10^{3 \cdot 23}$ -dimensional many-body wave function

- analytic solution does *not* exist
- numerical solutions only
- **1 Fe atom**, 10 data-points/dimension ... $10^{3 \cdot (26+1)}$ bytes
(1 DL-DVD $\sim 10^{10}$ bytes) ... 10^{71} DVDs ... "mission impossible"!
- such a solution is unnecessary and would also be a non-sense !
- further approximations are necessary !

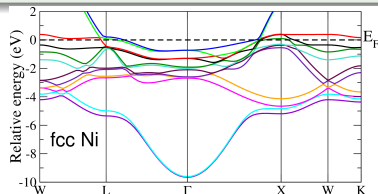
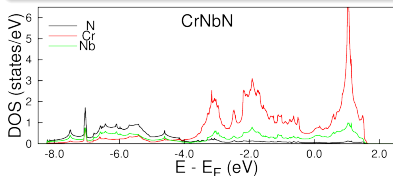
DFT calculation - Simple Example



Quantities from *Ab initio* calculations

Selected list of quantities

- Density of states, bandstructure, Fermi level, band gaps
- Structural stability of phases
- Magnetic order, non-magnetic, ferro-, antiferro-, spin-spirals, transition temperatures
- Elasticity, is it brittle or ductile? Young moduli, strength
- Defects, lattice vibrations, thermal expansion
- Optical and magneto-optical properties Kerr, Faraday, Voigt x-ray absorption
- Relativistic effects, superconductivity



The principle of XMLD

Microscopic description:

- XMCD ($\approx \mathbf{M}$) probes FM ordering only, sum rules gives us orbital and spin moment^a
- XMLD ($\approx \mathbf{M}^2$) probes FM and AFM ordering, anisotropy^b
- XMLD = $\frac{T_{\parallel} - T_{\perp}}{T_{\parallel} + T_{\perp}} = \frac{R(H_{\parallel}) - R(H_{\perp})}{R(H_{\parallel}) + R(H_{\perp})} \approx \frac{\omega d}{c} \text{Im}[n_{\parallel} - n_{\perp}] \approx$
 $\frac{\omega d}{2cn} \text{Im}[\varepsilon_{\parallel} - \varepsilon_{\perp} - \frac{\varepsilon_{od}^2}{\varepsilon_{\perp}}]$
- $\varepsilon_{\perp} = \frac{1}{2}(\varepsilon_{+} + \varepsilon_{-})$ then XMLD $\approx \varepsilon_{\parallel} - \frac{1}{2}(\varepsilon_{+} + \varepsilon_{-})$, high symmetry directions (4-fold) symmetry $\varepsilon_{\pm} = \varepsilon_{xx} \pm i\varepsilon_{xy}$
- XMLD in reflection for p -polarized light (not-normalized):

$$\text{XMLD} = R_{pp}(\mathbf{M}_x) + R_{pp}(-\mathbf{M}_x) - R_{pp}(\mathbf{M}_y) - R_{pp}(-\mathbf{M}_y)$$

This cancels the contributions linear in magnetization.

^aThole et al, PRL 68, 1943 (1992), Carra et al, PRL 70,694 (1993)

^bv.d. Laan PRL (1999), H. Kuneš and P. M. Oppeneer, PRB (2003)

For the 3p edges:

- Relativistic (spin-orbit) SCF cycle (WIEN2k)
- Routines for
 - ① Linear response theory - Kubo formula tensor $\varepsilon_{ij}(\omega)$
 - ② Momentum matrix elements ($M_i = \langle n' \vec{k} | \vec{p} \cdot \vec{e} | n \vec{k} \rangle^2$), *i.e.* among all bands and over all k-points
 - ③ Joint density of states - absorptive part of ε_{ij}
 - ④ In-house Kramers-Krönig relations \rightarrow complex ε_{ij} tensor elements
- XAS, XMCD, XMLD, etc. based on Yeh's 4-vector formalism^a
 - ① Optical response of multilayered structure for various polarization vectors (\mathbf{E}), incidence wave vector (\mathbf{k}) with respect to magnetization direction (\mathbf{M}) and crystal axis
 - ② Obtaining energy dependence of reflectivity (R_{pp} , R_{sp} , R_{ps} , R_{ss}) and transmittivity (T_{pp} , T_{sp} , T_{ps} , T_{ss}) for \rightarrow LMOKE, TMOKE, QMOKE, MLD, Schäffer-Hubert (Voigt) effect etc. for X-ray as well as for valence states.
 - ③ including grazing incidence

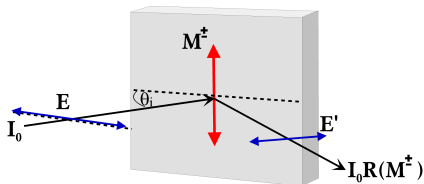
^aYeh (1980), Višňovský (1991)

Experimental setup & calculations on equal footing

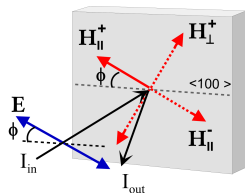
3nmAu/x nm Fe/150nm Ag/GaAs

$$A_{\text{TMOKE}} = \frac{R_{\text{pp}}(M_+) - R_{\text{pp}}(M_-)}{R_{\text{pp}}(M_+) + R_{\text{pp}}(M_-)},$$

$$A_{\text{XMLD}} = \frac{R(H_{\parallel}) - R(H_{\perp})}{R(H_{\parallel}) + R(H_{\perp})}, \quad R(H_{\parallel,\perp}) = \frac{R(H_{\parallel,\perp}^+) + R(H_{\parallel,\perp}^-)}{2}$$

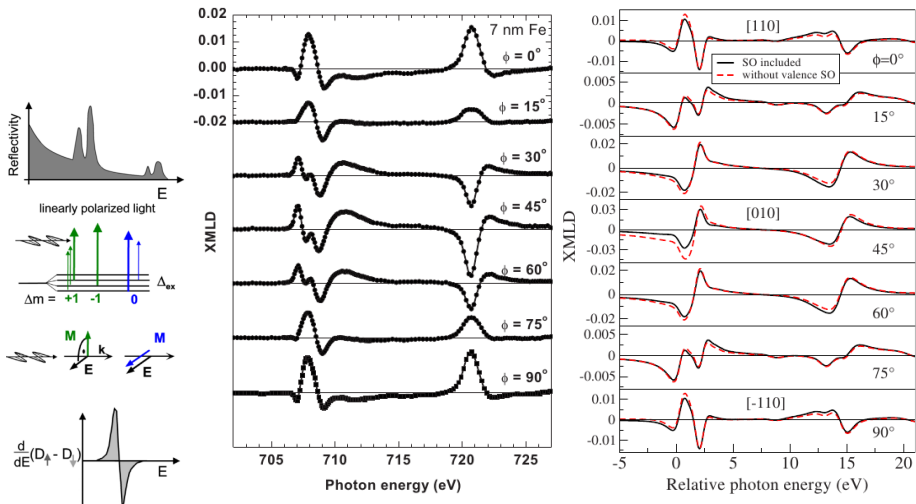


TMOKE setup



XMLD setup

XMLD at 2p edges in transmission - angular dependence



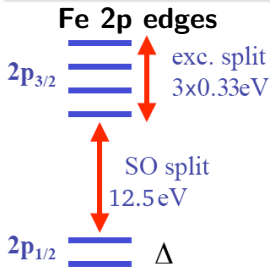
Nolting et al. PRB (2010), XMLD for high symmetry crystal directions at 2p edges,

theory by Kuneš and Oppeneer PRB (2003)

2p vs. 3p edges

2p edges: larger energy separation, no \mathbf{k} -dependence

3p edges: smaller energy separation \approx hybridization, \mathbf{k} -dependence

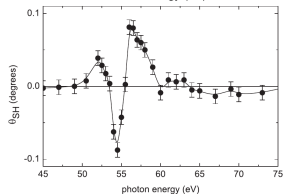
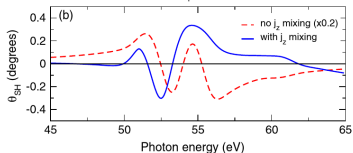
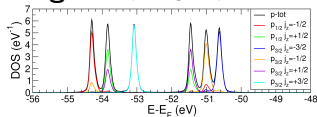


MO \approx M² (XMLD in reflection, Schaeffer-Hubert effect, etc.)

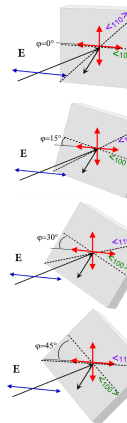
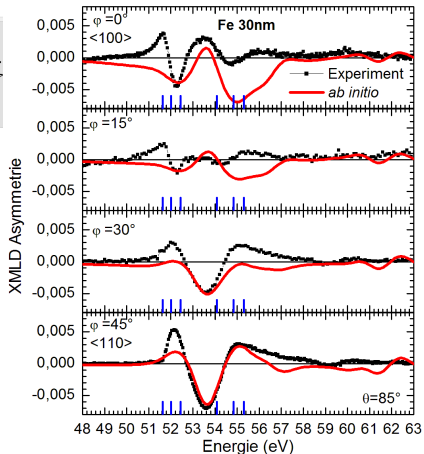
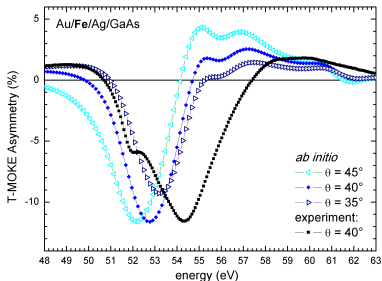
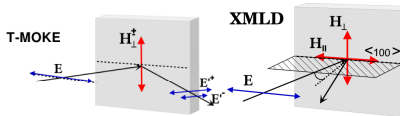
Valencia et al., Phys. Rev. B **104**, 187401

(2010)

3p edges J_z splitting of 3p-states in bcc Fe



TMOKE and XMLD of 3p Fe edge: Exp. vs. Theory

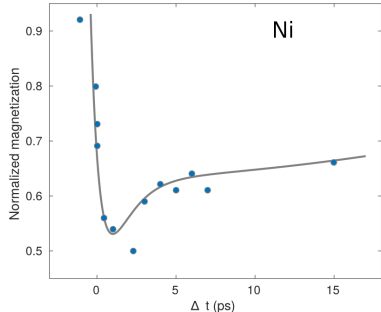
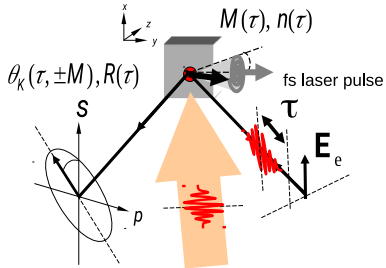


MO 3p edge is much susceptible to interference effects

Acta Phys. Pol. A **127**, 466 (2015), IEEE Trans.Magn., **50**, 2104704 (2014) J. Appl. Phys. **115**, 17E132 (2014), Phys. Rev. B **89**, 140404(R) (2014).

Start of the ultrafast demagnetization

Experiment opening the field of ultrafast demagnetization



Beaurepaire et al.(1996), magnetization quenched under 250fs.

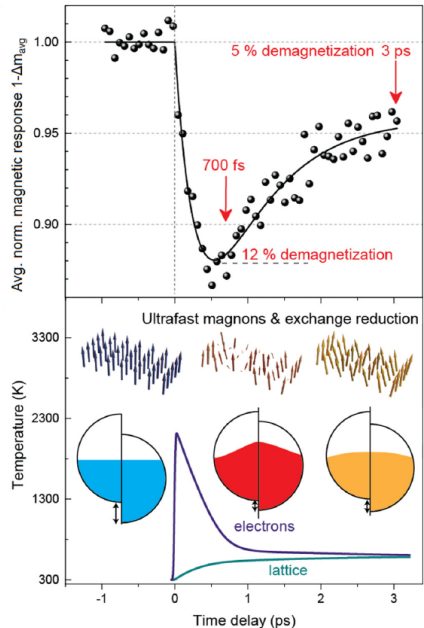
Theories of ultrafast demagnetization

- Elliott-Yafet el.-ph. scattering, Koopmans et al., PRL 95 (2005)
- Spin-orbital interaction $S \rightarrow L$, Zhang & Hübner, PRL 85 (2000)
- Fast magnon excitation, Carpena et al., PRB 78 (2008)
- Coulomb-exchange e-e scattering, Krauss et al., PRB 80 (2009)
- Chem. potential adjustment, Mueller et al., PRL 111 (2013)
- Spin-orbit inter. + opt. excitation, Krieger et al., JCTC 11 (2015)
- Fast $S \rightarrow L \rightarrow$ lattice transfer, Töws Pastor, PRL 115 (2015)
- Superdiffusive spin transport, Battiato et al., PRL 105 (2010)

Ultrafast demagnetization in hcp Co

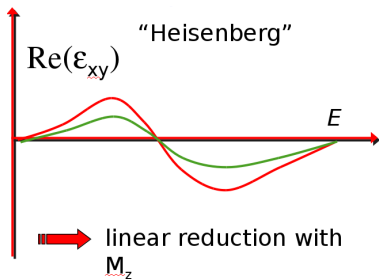
Three possibilities

- Magnetic moment disorder (transverse spin excitations)
 - Exchange splitting reduction (longitudinal spin excitations)
 - Electron temperature increase (“hot” electrons)
-
- 10nm Co/SiO₂
 - No superdiffusion
 - Electron thermalized at 400 fs
 - Measured at M-edges at ca 50eV



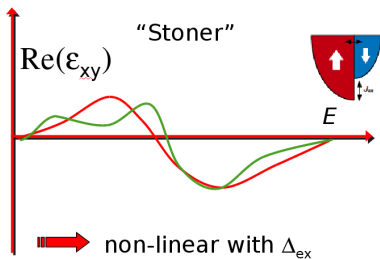
Magnon excitations vs. exchange excitations

Spin wave



Transverse excitation

Reduced exchange splitting



Longitudinal excitation

Oppeener et. al, Z. Phys. B 88, 309 (1992)

Can be these effects distinguished by T-MOKE?

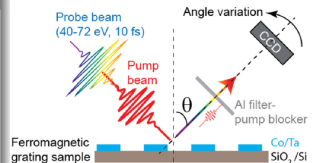
predicted by Erskine Stern, PRB 12, 5016 (1972)

- Measure whole energy and angle dependent T-MOKE asymmetry at 700 fs and 3 ps and determine change in $A(t)$ wrt $A(t=0)$, where A is

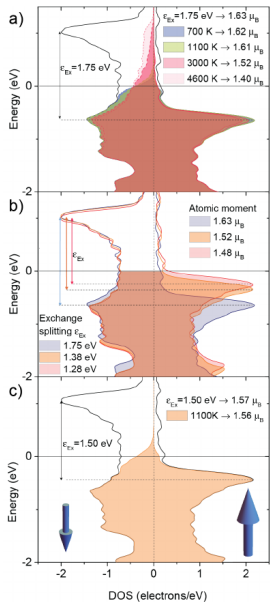
$$A(n, \theta, \varepsilon_{x,y}(t)) = \frac{I_p^+ - I_p^-}{I_p^+ + I_p^-} \approx 2\text{Re} \frac{\sin 2\theta \varepsilon_{x,y}(t)}{n^4 \cos^2 \theta - n^2 + \sin^2 \theta}$$

Compute the $\varepsilon_{x,y}$ from first principles (ab initio) for number of cases:

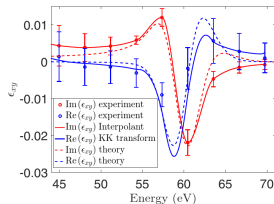
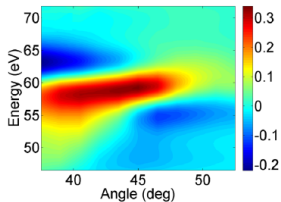
- magnon excitations
- reduced exchange splitting
- increased electron temperature
- construct $A(t)$ wrt $A(t=0)$, i.e. fitting with experiment



DOS for various excitations by relativistic code WIEN2k

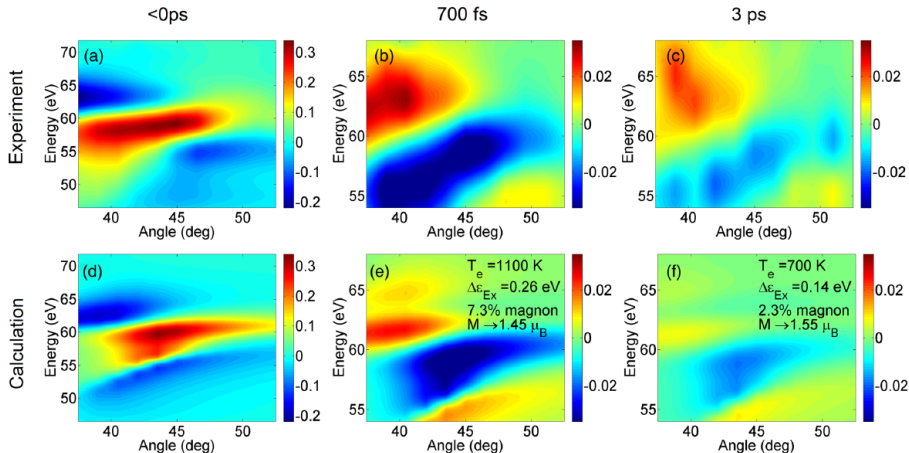


- a) hot electrons, FD multiplication
- b) reduce exchange, fixed spin moment calcs.
- c) combination of both

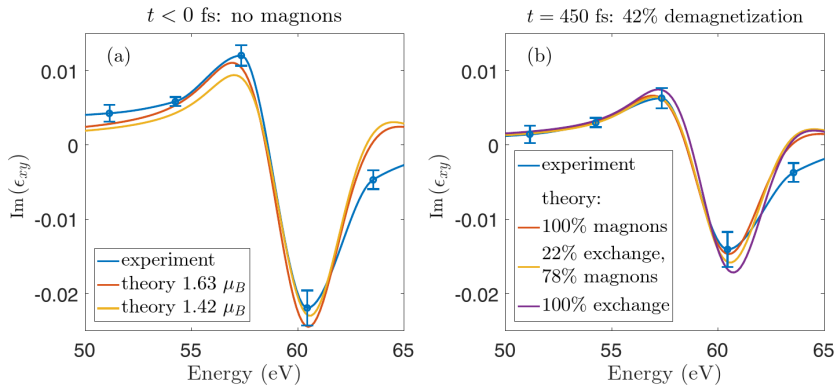


Experiment vs. theory - hcp Co

$$A(n, \theta, \varepsilon_{x,y}(t)) = \frac{I_p^+ - I_p^-}{I_p^+ + I_p^-} \approx 2\text{Re} \frac{\sin 2\theta \varepsilon_{x,y}(t)}{n^4 \cos^2 \theta - n^2 + \sin^2 \theta}$$

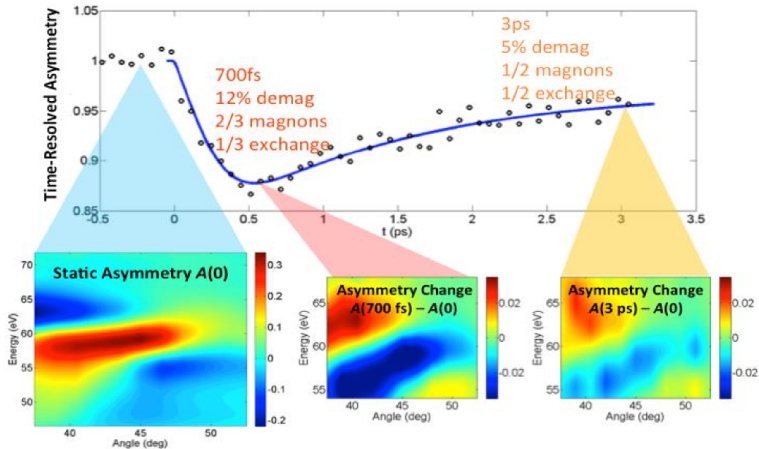


Least square fits to the whole spectrum, contribution from spin flips (exchange reduction) and magnons



Turgut et al., Phys. Rev. B **94**, 220408(R) (2016), Zusin et al. Phys. Rev. B **97**, 024433 (2018)

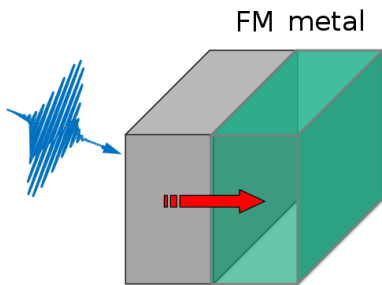
Causes of MOKE time evolution



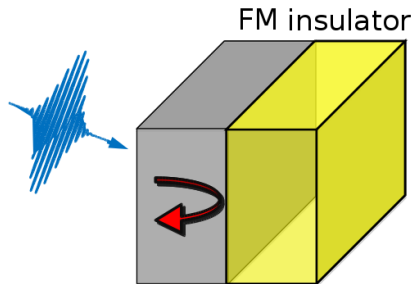
- surprisingly small contribution from spin flips (exchange splitting reduction)
- larger contribution by fast magnon excitations

Turgut et al., Phys. Rev. B **94**, 220408(R) (2016), Zusin et al. Phys. Rev. B **97**, 024433 (2018)

Emerging picture - Ultrafast demag. depends on substrate



- strong and fast demagnetization
- superdiffusive spins currents contribute predominantly



- smaller and slower demagnetization
- superdiffusive currents not possible
- due to spin wave generation and el-phonon spin flips

MO effects in transition metals

- MO at 3p edges are accurately calculated wrt experiment
- Linear/Quadratic MO effects at 3p edges are one order of magnitude lower than XMLD at 2p edges
- MO asymmetry measurement allows us to disentangle the different contributions to demagnetization
- Both transverse and longitudinal excitations found to contribute, at 700fs $\Delta m_t / \Delta m_l > 2$, later the contribution from longitudinal component is growing
- Magnon/exchange splitting is not growing from fs into ps time
- Ultrafast demagnetization of hcp Co on insulating (SiO_2) substrate by reduced exchange splitting

Thank you for your attention !

One post-doc and Ph.D. position open at IT4I, VSB - Tech. Univ.
Ostrava, Czech Republic

dominik.legut@vsb.cz