

USING SYMMETRIES

TO PREPARE GROUND STATES ON QUANTUM COMPUTERS

HPCSE 2026 20. May 2026

IVANA MIHÁLIKOVÁ, JOSEPH CARLSON, DUFF NEILL, IONEL STETCU

THE CHALLENGE: EXPONENTIAL HILBERT SPACE

N spin-1/2 particles → Hilbert space dimension 2^N

2^{12}

12 spins
4 096 states

2^{30}

30 spins
1 073 741 824 states

2^{50}

50 spins
1 125 899 906 842 624 states

Classical computers hit a wall. Quantum computers can help, and they help much more when they are guided to the relevant part of this space.

KEY QUESTION: HOW DO WE FOCUS THE ALGORITHM ON THE PART OF THE HILBERT SPACE THAT ACTUALLY MATTERS?

THE CHALLENGE: EXPONENTIAL HILBERT SPACE

N spin-1/2 particles \rightarrow Hilbert space dimension 2^N

2^{12}

12 spins
4 096 states

2^{30}

30 spins
1 073 741 824 states

2^{50}

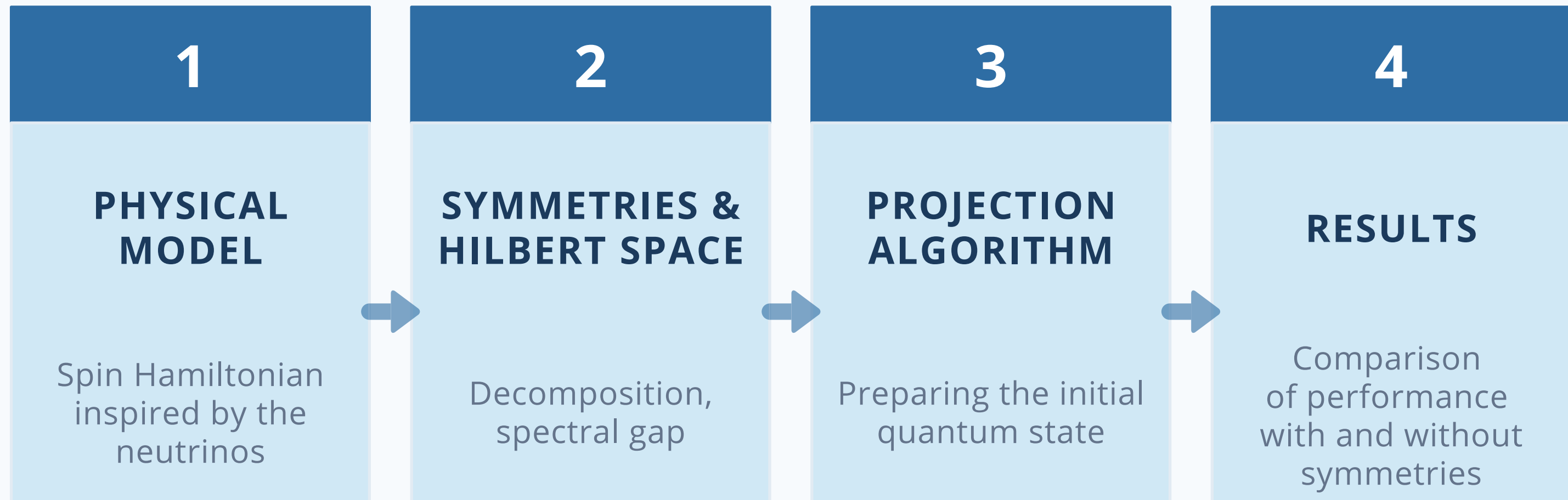
50 spins
1 125 899 906 842 624 states

Classical computers hit a wall. Quantum computers can help, and they help much more when they are guided to the relevant part of this space.

KEY QUESTION: HOW DO WE FOCUS THE ALGORITHM ON THE PART OF THE HILBERT SPACE THAT ACTUALLY MATTERS?

Symmetries reduce the relevant Hilbert space, increase the spectral gap, and give quantum algorithms a cleaner problem to solve, leading to faster and more reliable convergence.

OVERVIEW



1

PHYSICAL MODEL

Spin Hamiltonian inspired by the neutrinos

THE NEUTRINO-INSPIRED MODEL

All-to-All Random Couplings

$$H = \sum_{i < j} (1 - \cos \theta_{ij}) \sigma_i \cdot \sigma_j \quad (\text{all-to-all connections})$$

12 spins on a **4×3 lattice** with all-to-all random interactions

Full Hilbert space: $2^{12} = \mathbf{4096}$ dimensions

Coupling coefficient (1 - cos θ_{ij}): cos θ_{ij} is randomly chosen from the range [-1, 1]

Symmetries:

- total spin J,
- z-projection J_z

2

SYMMETRIES AND HILBERT SPACE

Decomposition, spectral gap

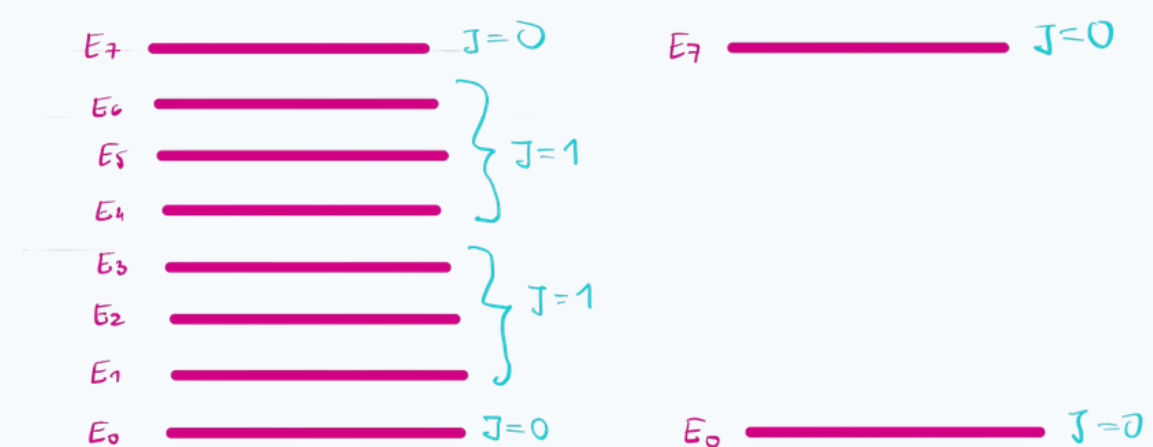
SYMMETRY AS HILBERT SPACE DECOMPOSITION

When H commutes with a symmetry operator \rightarrow Hilbert space decomposes into independent subspaces, each labelled by a quantum number (e.g. total spin J).

- States in different subspaces never mix under time evolution.
- If the ground state has $J = 0$, we restrict the entire algorithm to that subspace and ignore everything else.

KEY INSIGHT: WE DO NOT NEED TO EXPLORE THE FULL HILBERT SPACE — SYMMETRY TELLS US EXACTLY WHERE TO LOOK.

SYMMETRY AS HILBERT SPACE DECOMPOSITION



Symmetry gives us: smaller space, larger gap and more focused algorithm.

3

INITIAL STATES AND PROJECTION ALGORITHM

Preparing the initial quantum state

STARTING POINT — PRODUCT STATE

- Simple product state: tensor product of single-qubit states, easy to prepare with single-qubit gates only
- **Neutrino model:** flavor-momentum locked state; each spin aligned to its momentum vector (semiclassical ground state for large N)

$$|\psi_i\rangle = \bigotimes_{i=1}^N \exp\left(i\pi k_i \cdot \frac{\sigma_i}{2}\right) |0\rangle_i,$$



- This state capture some physics but do not respect quantum symmetries

MODEL	INITIAL FIDELITY
Neutrino model	2.4%

$$\text{Fidelity} = |\langle \psi | \varphi \rangle|^2$$

ψ = prepared state φ = exact ground state

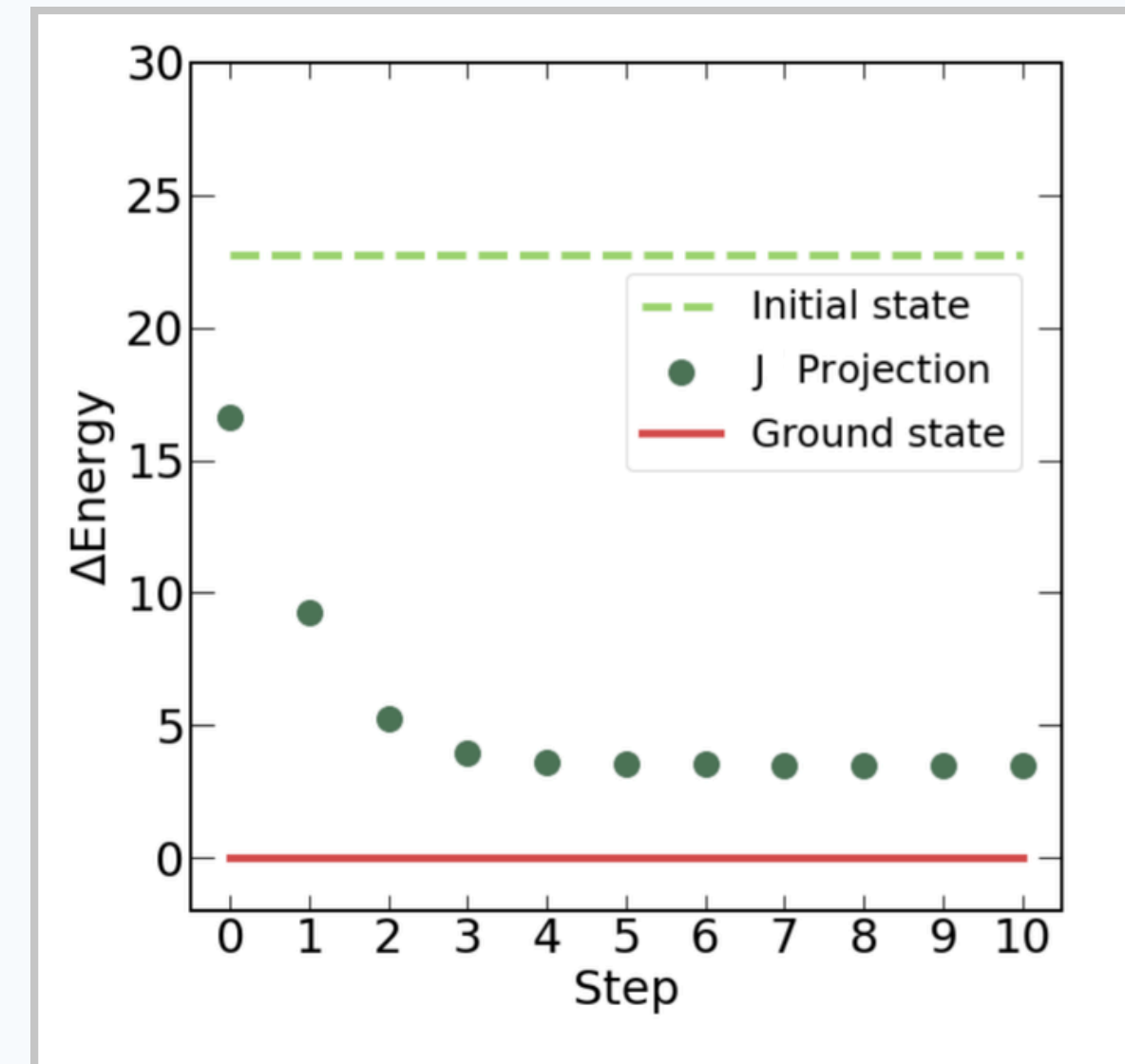
Low initial fidelity but a good warm start for the projection step.

THE PROJECTION ALGORITHM

Alternating projections onto $J_z = 0$ and $J_x = 0$, each followed by a measurement of the ancilla, iteratively filter out the $J > 0$ components. 11 steps are enough to reach 56.7% fidelity.

$$|\psi(t_i)\rangle = e^{-it_i J_z \otimes Y_a} |\psi\rangle \otimes |0\rangle_a$$

$$|\psi(t_i)\rangle = e^{-it_i J_x \otimes Y_a} |\psi\rangle \otimes |0\rangle_a$$



~80% of the way to exact ground state after 11 steps

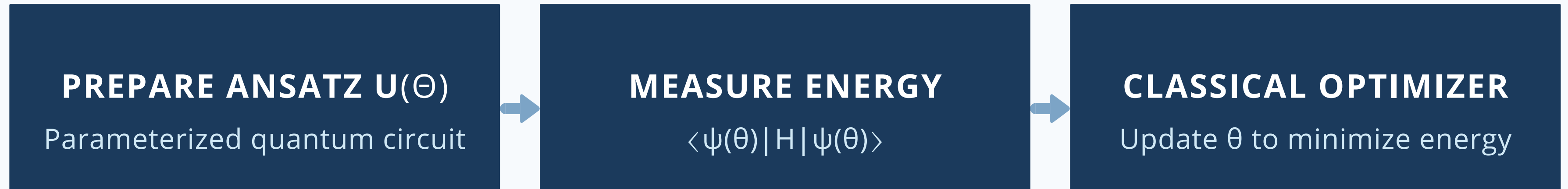
After projection: fidelity 2.4% \rightarrow 56.7%

4

RESULTS

Comparison of performance with and without symmetries

VARIATIONAL QUANTUM EIGENSOLVER



Repeat until convergence

- **Variational principle:** $\langle \psi | H | \psi \rangle \geq E_0$ for any trial state, minimizing energy gives best approximation
- **Main challenge:** barren plateaus; gradient becomes exponentially small, optimizer gets stuck
- **Solutions:** start close to the correct answer, reduce number of parameters
- Symmetry-projected initial state is an ideal warm start; fidelity already **57%**

VQE: WITHOUT VS. WITH SYMMETRY PROJECTION

Without Projection

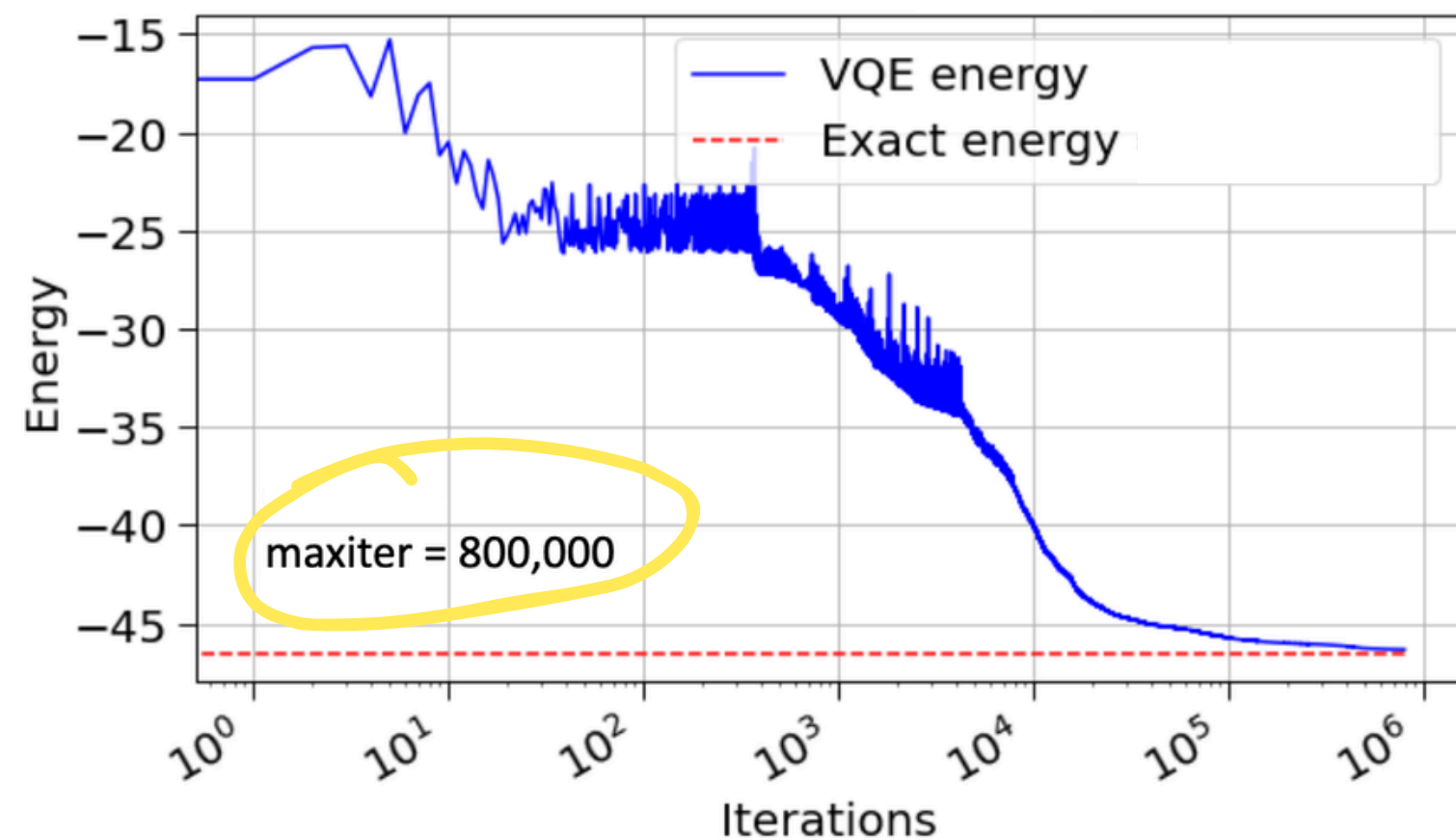
- Searches full 4096-dimensional Hilbert space
- No guidance on which subspace matters
- Prone to barren plateaus and local minima
- May converge to wrong symmetry sector
- Low final fidelity with exact ground state

With Projection

- Searches only the correct symmetry subspace
- Starts from high-fidelity projected state
- Better-behaved optimization landscape
- Swap ansatz preserves J and J_z throughout
- Fidelity: >98%

VQE: WITHOUT VS. WITH SYMMETRY PROJECTION

Without Projection

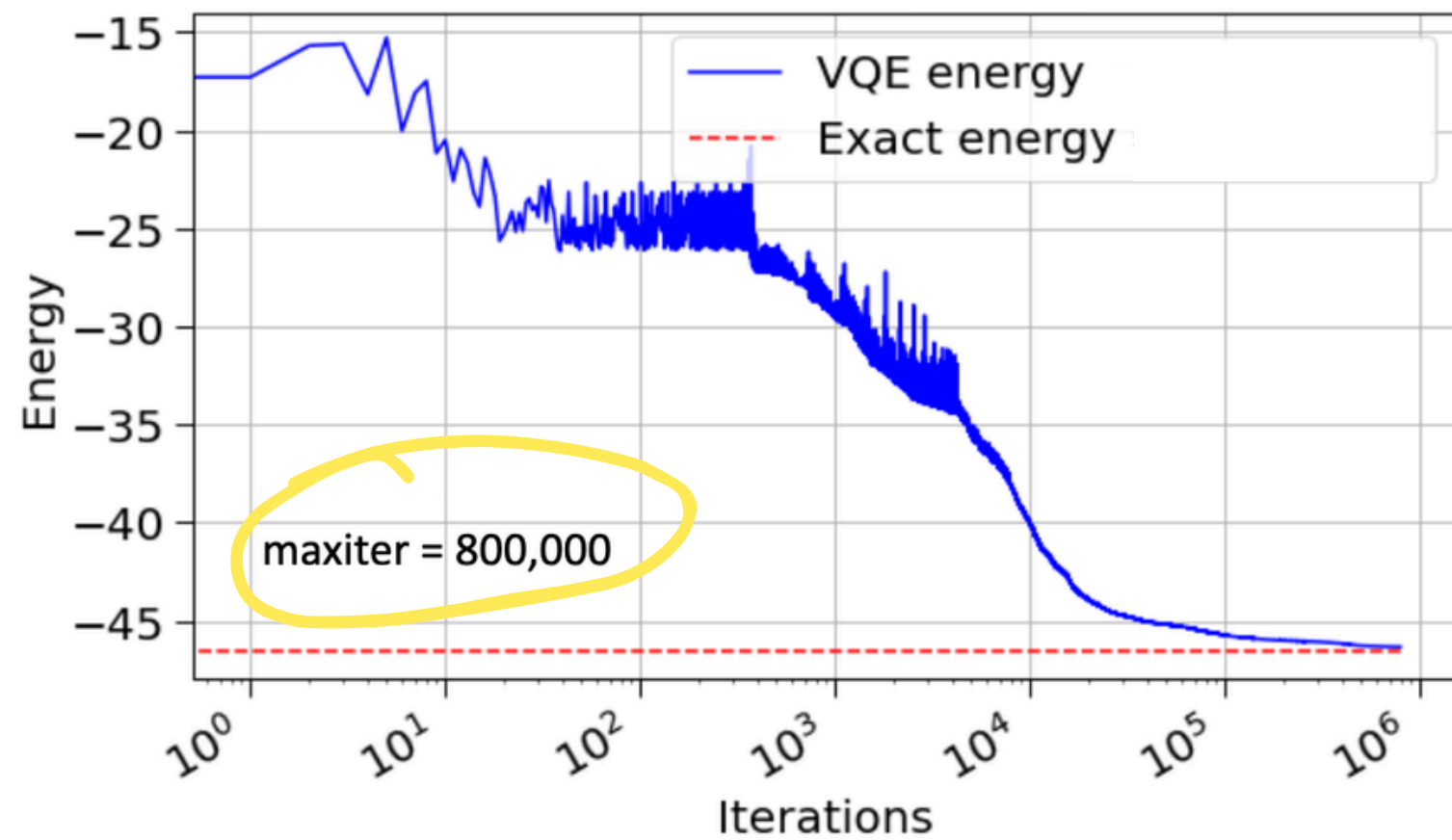


With Projection

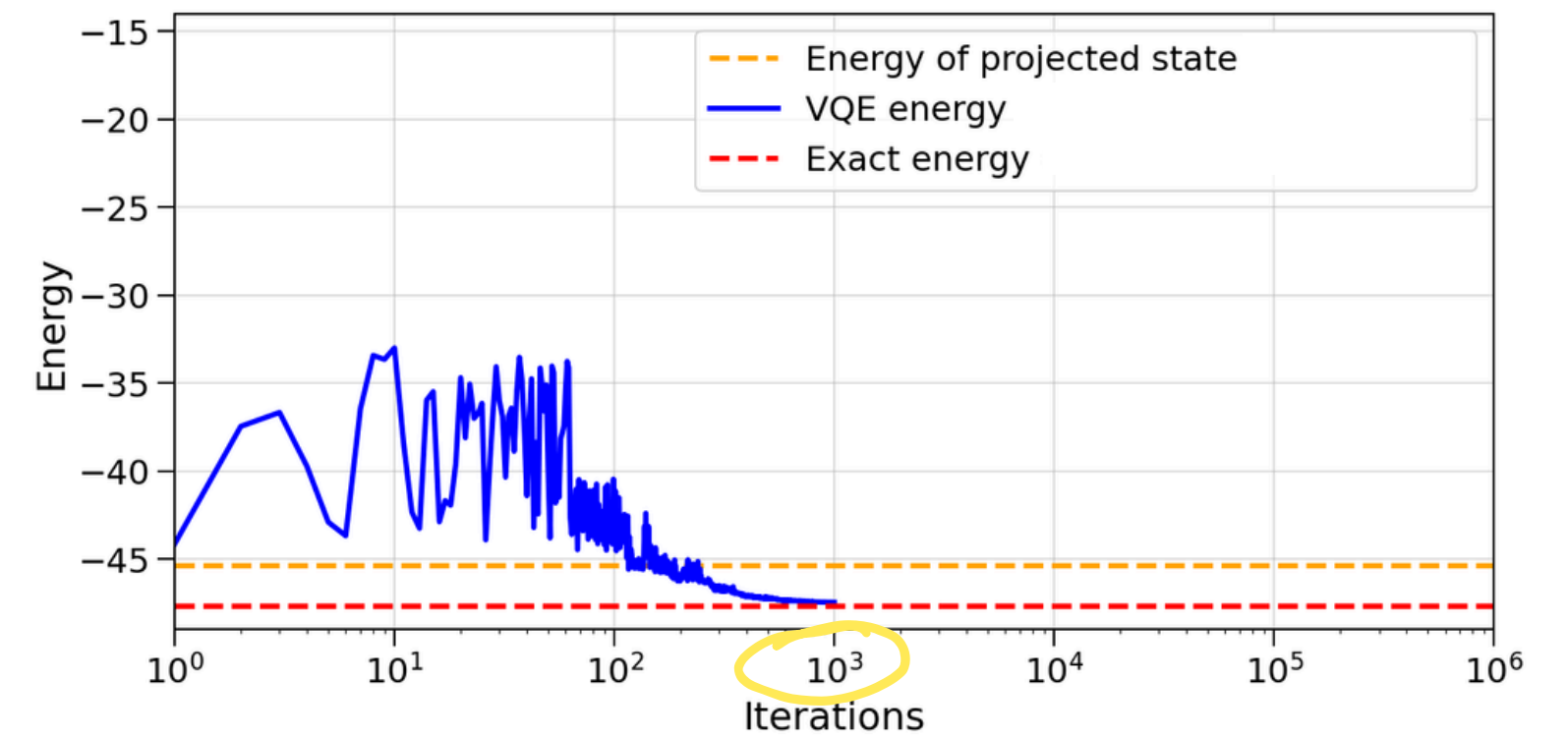
- Searches only the correct symmetry subspace
- Starts from high-fidelity projected state
- Better-behaved optimization landscape
- Swap ansatz preserves J and J_z throughout
- Fidelity: >98%

VQE: WITHOUT VS. WITH SYMMETRY PROJECTION

Without Projection



With Projection



CONCLUSIONS

arxiv: 2510.06702

STEP	FIDELITY
Initial product state	2.4%
J ² Projection (11 steps)	56.7%
VQE convergence	98.8%

Symmetry:

- Reduces the relevant Hilbert space from 4096 to 132 states
- Improves the initial state from 2.4% to 56.7% fidelity in just 11 steps
- Guides VQE to 98.8% fidelity with the exact ground state
- Makes the algorithm 800× faster

THANK YOU FOR YOUR ATTENTION!